

# Temperature and Emissivity of a Shocked Surface: A First Experiment

*P. Poulsen and D. E. Hare*

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# Temperature and Emissivity of a Shocked Surface: A First Experiment

Peter Poulsen and David E. Hare  
Lawrence Livermore National Laboratory

We have conducted an experiment in which the temperature and the wavelength dependent emissivity of a shocked surface has been measured. In the past, only the thermal emission from the shocked surface has been measured. The lack of knowledge of the emissivity as a function of wavelength leads to uncertainty in converting the measured emission spectrum into a surface temperature. We have developed a technique by which we are able to calculate both the emissivity of the shocked surface over a range of relevant wavelengths and the temperature of the surface. We use a multi-channel spectrometer in combination with a pulsed light source having a known spectrum of infrared radiation. Two separate techniques using a pulse of reflected radiation are employed and described. Both give the same result: An initially polished molybdenum surface that is shocked and partially released has a temperature of 1040 degrees Kelvin and a wavelength ( $\lambda$ ) dependent emissivity of 0.16 ( $\lambda=1.2\mu\text{m}$ ), 0.10 ( $\lambda =1.6 \mu\text{m}$ ), and 0.20 ( $\lambda =2.3 \mu\text{m}$ ).

## The problem with measuring temperature

The measurement of the temperature of a material is essential for understanding a wide range of materials processes (e. g. shock dynamics, foundry operations, and wafer manufacture) where knowledge of the state of the material is important. Inferring the temperature from a measurement of radiation is difficult, however. An experiment in which the thermal emission from a shocked material was successfully measured is described in reference 1. In that experiment, the emitted power from a warm surface is measured over a range of wavelengths using a multi-channel spectrometer. Each channel of the spectrometer passes the radiation from a selected wavelength interval into a detector. The spectrometers typically have anywhere from three to six channels, and not all channels may have enough signal to contribute to the measurement under any given condition. The difficulty in the measurement lies in the relation of the emitted power in each channel to the temperature of the surface. The emitted power is determined not only by the temperature of the surface, but also by the emissivity of the surface. The emissivity is not a fixed number for any real surface, but varies both with angle of observation and with wavelength. Thus the temperature cannot be calculated from the emitted power spectrum without detailed knowledge of the emissivity. Approximate temperatures may be calculated by various assumptions, such as assuming a fixed number for the emissivity (gray body), assuming the emissivity in adjacent channels is the same, causing the relative emissivity to drop out of the relations giving the power ratio, and by assuming a lower limit on the emissivity. The latter method will yield a

range of temperatures and emissivities in each channel which are consistent with the measurement.

It is possible, as shown below, to solve explicitly for the emissivity in each channel wavelength interval as well as for the temperature of the surface. The additional data that make this possible can be either a measurement of the spectrum of light reflected from the surface from a source of known spectral content or a measurement of the reflected power from both the shocked surface and a calibrated surface. Note that it is not necessary to know the absolute values of the flux, either reflected or emitted. This is a great advantage because the exact geometry of the configuration may be rapidly changing or not known at all.

This note discusses the relations that are used and needed for the calculations as well as a little of the theoretical basis. The relations assume that we have measured the reflected and emitted power from a number of wavelength bands of the surface. We have previously conducted a demonstration experiment (reference 2, 3) in which a stationary, steady state surface was used as the source of emitted thermal radiation. We also acquired a fast source (risetime  $< 0.5 \mu\text{sec.}$ ) of light that extended into the infrared regime, and used that to acquire the reflectivity data. We know that we can measure the emitted light in 0.1 - 0.5 microseconds from previously published experiments (reference 1). We also know from the same experiments that we can maintain a steady state condition on a shocked surface for approximately one microsecond. We therefore have a tool that is capable of measuring both reflected and emitted light in less than a microsecond and which is therefore applicable to the determination of the temperature and emissivities of a shocked surface. This note discusses the results of the gas gun experiment in which we were able to approximate the required conditions and obtain the surface temperature of a shocked surface.

## **Discussion of the methods**

We will discuss two related but separate methods for obtaining the emissivities and the surface temperature. As noted above, both methods use a pulse of light incident onto the surface to acquire information about the reflectivity of the surface. In the first method, here named Method I for lack of a more imaginative name, we only need to know the spectrum of the incident light as measured by the spectrometer. The relations do not require knowledge of the geometry and absolute determination of the emitted and the reflected power as long as the light is measured at the same point where the incident light originates, a condition which can be reasonably well approximated in an experiment. It is very useful to have a diagnostic which can determine the temperature of a surface without knowledge of the exact geometry relations. The method does have limitations which originate in the sensitivity of the solution under certain conditions. These conditions, related to the magnitude and variation of the emissivity with wavelength, are discussed below.

In the second method, inventively named Method II, we substitute a surface with a known reflectivity for the surface of interest (the one that will be shocked). We then measure the light reflected from the substituted surface, replace it with the surface of interest, and measure the reflected light from that surface; in this experiment, the measurement takes place after the surface has been shocked. The reflectivity

and the emissivity as a function of wavelength of the shocked surface is calculated by comparing the two measurements of reflected power. The temperature is then computed by finding the temperature for which the ratio of the measured power to computed power in each wavelength interval is the same. This method is generally applicable with no restrictions on the value of the emissivity other than measurement accuracy. The method does require that the viewing geometry is effectively the same for both the calibration surface and the shocked surface. The shocked surface can move (most do!), but the viewing conditions must be such that the power emitted and reflected is independent of distance. This condition is usually easily satisfied since the surface to be viewed must have a uniform temperature and emissivity within the angle of acceptance of the diagnostic. The disadvantage of this method, of course, is that it may not always be possible to make the substitution of the surface to be measured with the calibration surface.

### **The equations for method I**

This method uses the relative magnitude of the emitted light in each wavelength channel and the incident light spectrum measured in the same channels to compute the channel emissivities and the surface temperature.

The relations for the emitted light are:

$$S_i = \epsilon_i F_i (T) K C_i$$

The relations for the reflected light are:

$$R_i = ( 1 - \epsilon_i ) I_i H C_i$$

Note that it has been assumed that the reflectivity is unity minus the emissivity for each wavelength. There are conditions relating to the angular dependence of the emissivity that must be satisfied for this relation to be applicable. They are discussed below.

The definitions of the symbols are as follows:

H and K are geometry factors. They can be made to be, and are assumed to be, independent of the wavelength channel.

i – denotes channel i

$C_i$  – sensor calibration factor for channel i

$F_i (T)$  – black body power in channel i

$I_i$  – power from flash lamp source in channel i

$\epsilon_i$  –emissivity of surface in the wavelength interval of channel i

$S_i$  – emission signal in channel i

$R_i$  – signal from reflected flash lamp in channel i

It is highly desirable to eliminate the geometry factors since we'd like to work with relative values only; the use of relative values, i.e. the spectrum instead of absolute flux, eliminates serious problems with absolute calibration and the instantaneous measurement of geometry. Hence we choose an arbitrary channel “n” and normalize the equations. Typically the chosen channel is the one with good signal resolution for both the emitted and the reflected signal. The equations become:

$$A_i = e_i B_i (T)$$

$$D_i = E_i (1 - e_i \epsilon_n) / (1 - \epsilon_n)$$

with

$$e_i = \epsilon_i / \epsilon_n$$

$$A_i = S_i C_n / S_n C_i$$

$$B_i = F_i / F_n$$

$$D_i = R_i C_n / R_n C_i$$

$$E_i = I_i / I_n$$

We know the values of  $A_i$ ,  $D_i$ ,  $E_i$ , and the functional dependence of  $B_i$  on  $T$ . This set of equations allows the solution of  $e_i$ ,  $\epsilon_n$ , and  $T$ . Every set of three equations yield a temperature and a set of emissivities if the reflectivity relations are used. In the gray body case, the reflectivity relations contribute no information, the relative emissivities are unity, and any two of the equations yield a temperature. The spread of calculated temperatures is a measure of the consistency and accuracy of the data.

The solution to the above equations can under some circumstances be extremely sensitive to variations in the input data and the experimental parameters. It is worth a brief discussion here since the regions of sensitivity in the equations affect the way the experiment should be done. A more detailed discussion of this subject will be published at a later date.

There are two conditions that cause difficulty: When the emissivities are very nearly equal in each channel, and when they are very small. In the first case the solution is clear: The reflectivity measurement is not needed and the calculation of the relative emissions in pairs of channels should each give the same temperature for ideal data input. The second condition is more difficult experimentally. If the emissivities are all small so the reflectivities are near unity, but yet the ratios of the emissivities in pairwise channels

differ greatly from unity there is a serious problem that can only be solved by extreme accuracy in the data. The experimenter is advised to estimate the required accuracy prior to designing the experiment.

The emissivity of a surface is a function of both wavelength and the angle to the surface. The multichannel method described above approximates the wavelength dependence in a step-wise fashion; more channels result in greater accuracy as long as decreasing the filter bandwidth does not affect the signal to noise ratio. Otherwise the two factors must be traded off to optimize the diagnostic.

The angular dependence can be dealt with in either of two ways which render the geometry of emission and reflectivity effectively identical. The first way is to have the detector aperture and the source of pulsed light occupy the same point in space. That condition, of course, may not be perfectly satisfied, but closely spaced fiberoptics is good enough. The other way is to employ an integrating, perfectly reflecting enclosure. The latter has the advantage of increasing the effective emissivity and thus the emitted signal. The disadvantage is that it may be difficult to maintain uniform conditions for the shocked surface within the enclosure. Practically speaking, the interaction between the enclosure and the shocked surface may be a source of unwanted light.

It is noted that the equation set above has no constraints, i. e. there is nothing in the equations that force the emissivities to be positive and less than unity. Only good data will do that, and even then it is important to have an estimate of the errors.

## **The equations for Method II**

A measurement of the reflected power from a surface having a known reflectivity provides a basis for computing the reflectivity of the shocked surface. We use a gold coated surface which has a reflectivity near unity for the range of wavelengths of interest. The reflectivity of the surface  $r_i$  in the  $i$ 'th wavelength interval is calculated from the known reflectivity  $r_{ci}$  from the corresponding signals  $S_i$  and  $S_{ci}$  (subscript c denotes the calibration surface):

$$r_i = S_i / (S_{ci} / r_{ci})$$

The calibration factors cancel out in this calculation, so they are not needed. It is important, however, to maintain all conditions other than the surface identical for the two measurements.

Relating the emissivity to the reflectivity, we get

$$\epsilon_i = 1 - r_i$$

Knowing the emissivity, we can now calculate the ratio of the measured power to the calculated power. We normalize to the  $n$ 'th channel to eliminate geometry factors. The computed quantity is  $A_i / \epsilon_i B_i$  (T) using the above nomenclature. The solution point for the temperature occurs when this quantity has the same value for all channels.

## The Experiment

A sketch illustrating the experimental arrangement is shown in figure 1. A copper flyer driven by a gas gun to a velocity of 3 km/sec impacts a molybdenum plate backed by a lithium fluoride (LiF) window. The LiF window serves two purposes: It increases the pressure at the interface relative to a completely released condition, and it prevents the creation of ejected micro-particles at the surface. The difficulties associated with the LiF/metal interface are discussed in reference 1.

It is very important in any experiment in which a spectrum is measured to avoid any source of extraneous light. That is accomplished in this experiment by enclosing the LiF crystal with a non-reflecting surface. The only entry or egress of light is through the optical fibers. Both ends of the fibers must be carefully shielded from extraneous light, of course. As mentioned above, a fully reflecting surface may also be employed with care to integrate over angles and to increase the effective emissivity.

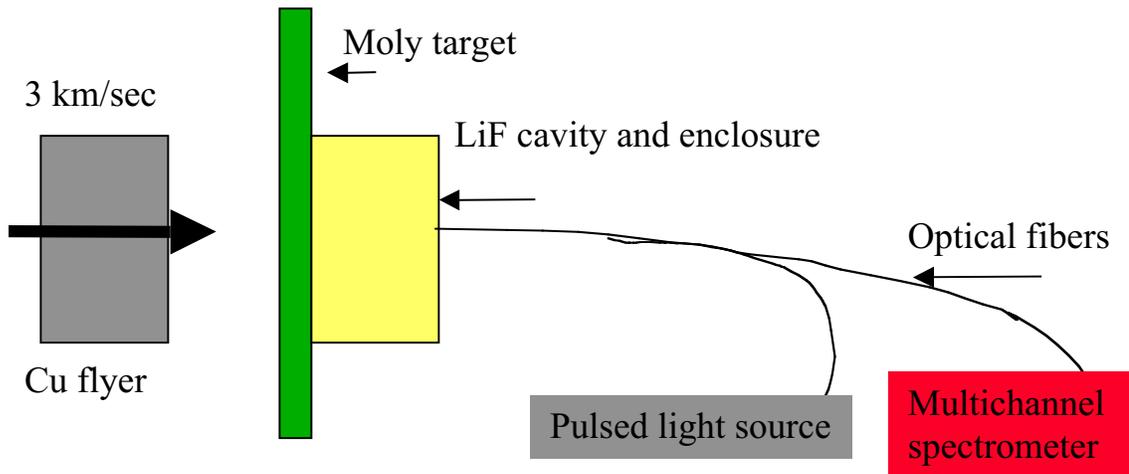


Figure 1. A schematic of the experiment

The idea of the experiment is to shock the moly plate and have the shock reach the interface between the moly and the LiF. The radiation from the shocked surface emits through the LiF and travels through the fiberoptics to the detectors where the signals are generated. Once a steady state has been reached at the surface (its duration is calculated to be approximately 1.0 microseconds), the source of pulsed light is fired and infrared radiation is reflected off the shocked surface; the sum of the emitted and the reflected radiation is measured by the detectors. In the experiment described below we fired the lamp a little too late because the flyer had a higher speed than expected (we expected 2.6 km/sec and measured 3.05 km/sec). The waveforms are illustrated in figure 2 below. The intent was to fire the lamp after the emitted signal had been measured, at a point corresponding approximately to  $-2.0$  microseconds and when the peak of the reflected signal was within the steady state region of the signal. We need to be able to subtract the

emitted signal from the total signal to obtain the reflected signal. We were still able to use the data due to additional knowledge about the waveform of the emitted light, as discussed below.

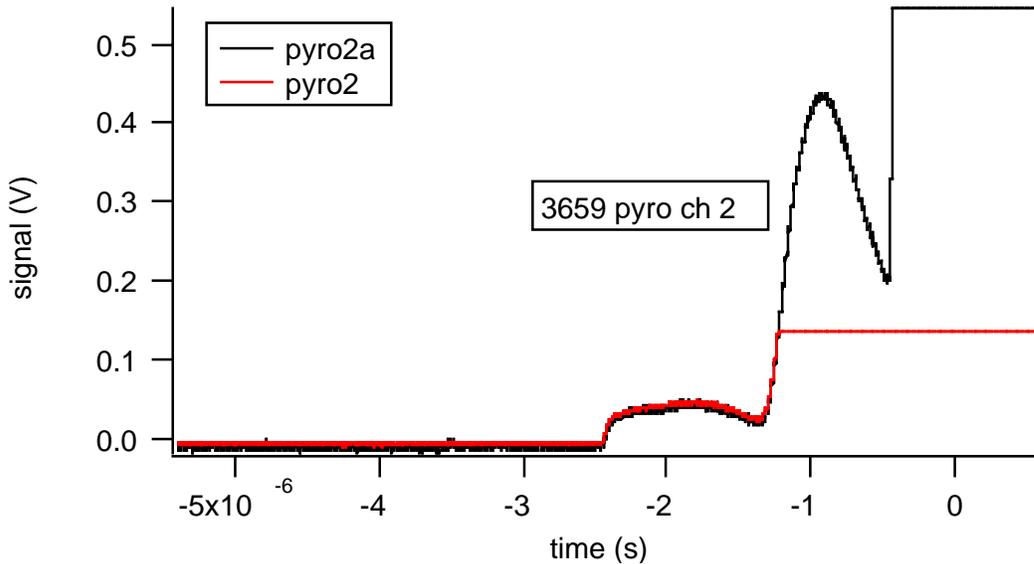


Figure 2. The waveform from the experimental channel 2 detector showing the rise in the signal due to the emitted radiation at approximately  $-2.4$  microseconds. The lamp is fired at approximately  $-1.3$  microseconds, later than desired as discussed in the text.

The target is shown in figure 3. The target is 6 mm plate of molybdenum with a polished surface, backed by the LiF crystal. The crystal is completely enclosed in a case which has an aperture for the fiberoptic coupling. The fiberoptic coupling is shown in figure 4.

These data were obtained with the LiF enclosure darkened by soot to minimize reflections. This minimized the measurement of radiation originating outside the cone of acceptance of the fiberoptics. It is important to have a surface of uniform temperature within the cone of acceptance of the fiberoptics and to have the geometry of emission the same as the geometry of reflectance. A region of non-uniform temperature is formed at the edge of the impact region due to release processes, and the contribution from this region to the total signal must not be significant if the measurement is to be interpreted by the above analysis methods.

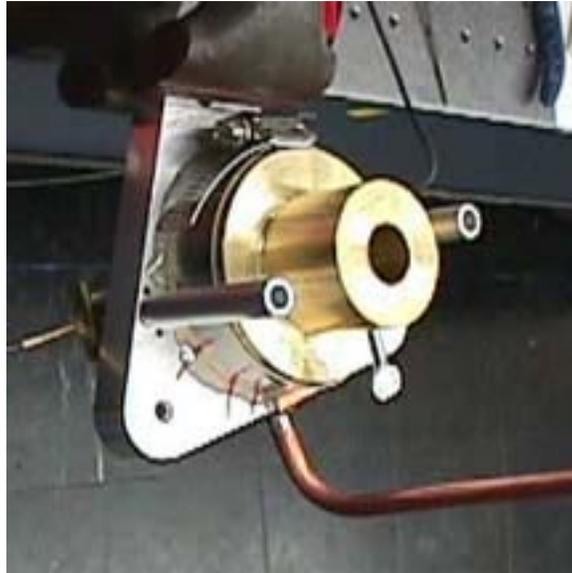


Figure 3. The moly target is in the plane of the triangular holder; the LiF crystal is covered by a light tight enclosure that holds the fiberoptic coupler.



Figure 4. The fiberoptic coupler is inserted into the enclosure. The coupler connects the fibers from the detector and from the lamp to the enclosure.

## The Detector

The detector has 6 channels, each selected for a particular wavelength region. Each channel has an IR detector, optics, a bandpass filter, and passes the light from the surface to the detector via optical fibers. It is similar to the detector described in reference 1 and is shown in figure 5. The configuration of the detector is shown in table 1.



Figure 5. The 6 channel infrared spectrometer

Table 1. Detector configuration – shot 3659

Pyro channel	Analysis channel	filter (um)	lenses	fiber	detector
ch1	NA	NA	NA	NA	NA
ch2	Ch 3	1.545 – 1.742 um	silicate glass	low OH FS	InSb
ch3	Ch 4	2.004 – 2.502 um	ZnSe	C2	InSb
ch4	Ch 2	1.130 – 1.348 um	Silicate glass	low OH FS	InSb
ch5	Ch 6	2.985 – 4.979 um	ZnSe	C2	InSb
ch6	Ch 5	2.503 – 3.490 um	ZnSe	C2	HgCdTe
lamp fiber (four)				C2	

## The Light Source

The source of pulsed light is a Hamamatsu xenon high pressure lamp with a sapphire window. It is not perfect for this application because it rises too slowly (about 0.5 microseconds to the peak) and because it has too little power in the 3 – 5 micron wavelength region. We have found it difficult to acquire a source of light in the infrared with the correct properties, and this source allowed us to use three channels and to show that the method works. This is an area that needs improvement and we have acquired other sources.

The temporal behavior of the pulse in one of the channels is shown in figure 6. It is desirable to have a really fast risetime because that allows the use of thinner targets and thus minimizes the effect of edge release on the data.

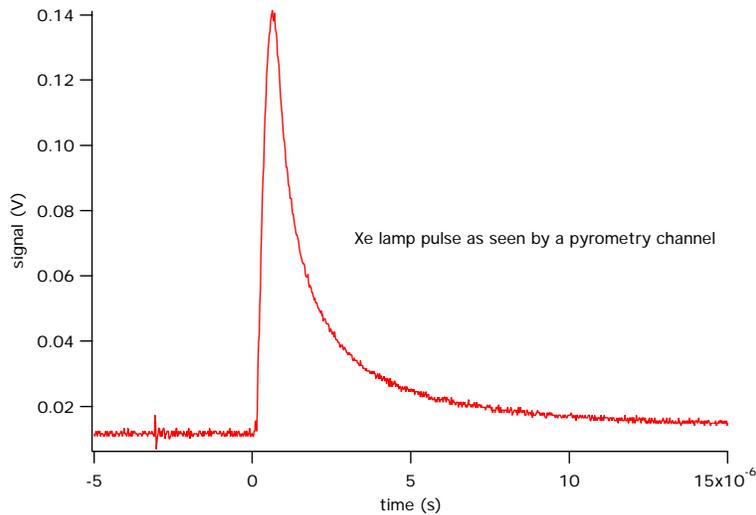


Figure 6. The Xenon lamp pulse signal as recorded in one of the pyrometer channels. It has a 0.5  $\mu$ second risetime; the risetime can be shortened further by changing the capacitor arrangement

As noted above, the power in the spectrum drops off at longer wavelengths. Raw data (uncalibrated) is shown in figure 7 as an illustration. We were unable to use the two large wavelength channels since the reflected power in these channels was insignificant relative to the emitted power.

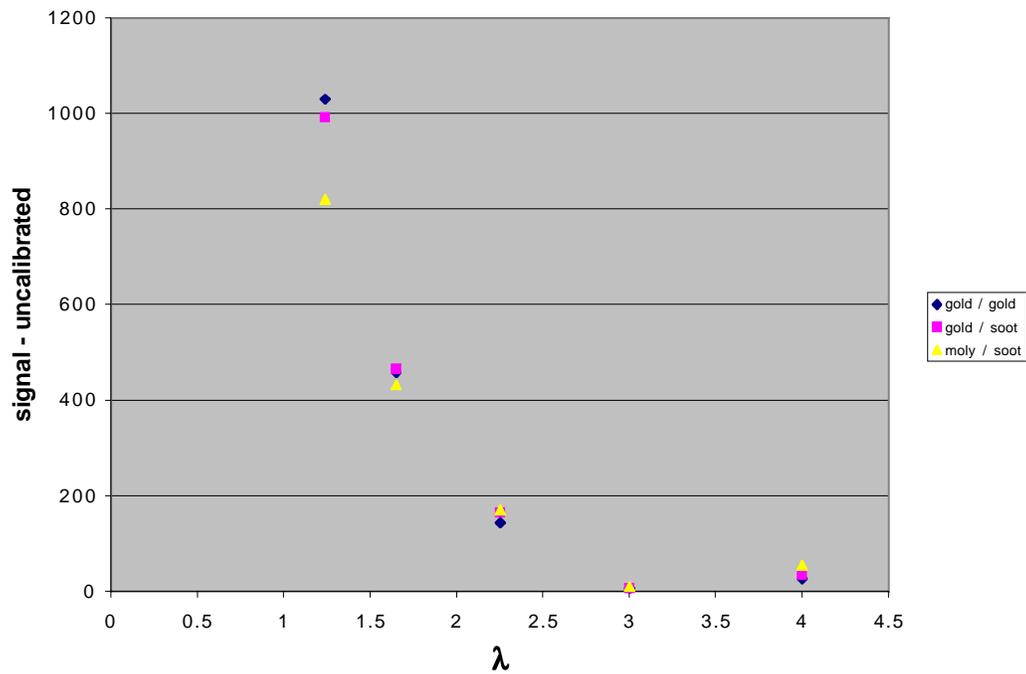


Figure 7. The uncalibrated power spectrum of the pulsed Xenon lamp; the blue and purple points are power reflected from a gold coated surface and the yellow points are power reflected from a polished molybdenum surface.

### The raw data

It is illustrative to include the raw data in this report; it is shown in figure 8. Note that the channel numbers used in the analysis, which increase with wavelength, are not the same as the detector channel numbers. The latter are only used to identify the channels.

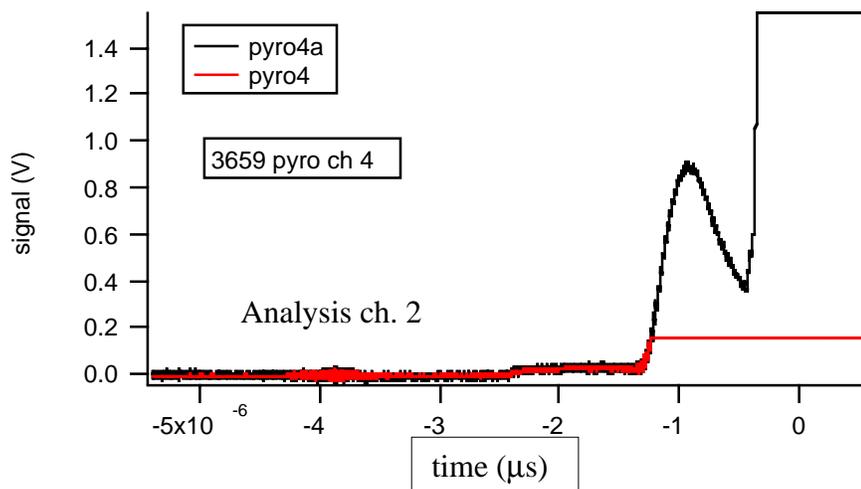


Figure 8a. Signal from analysis channel 2.

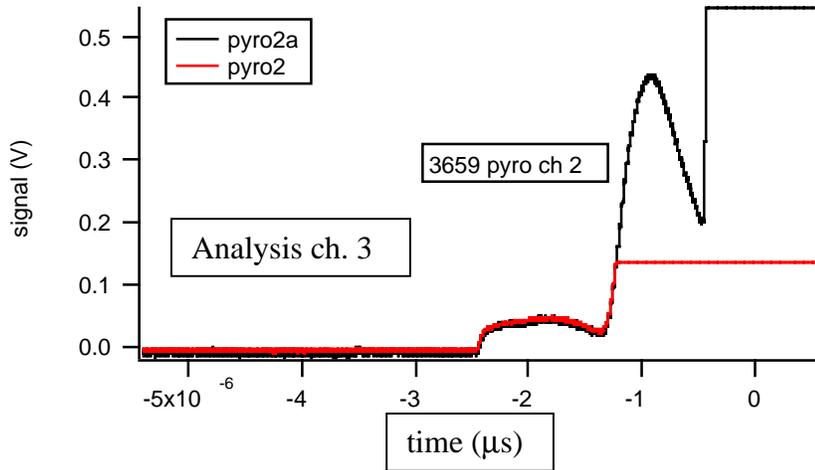


Figure 8b. Signal from analysis channel 3.

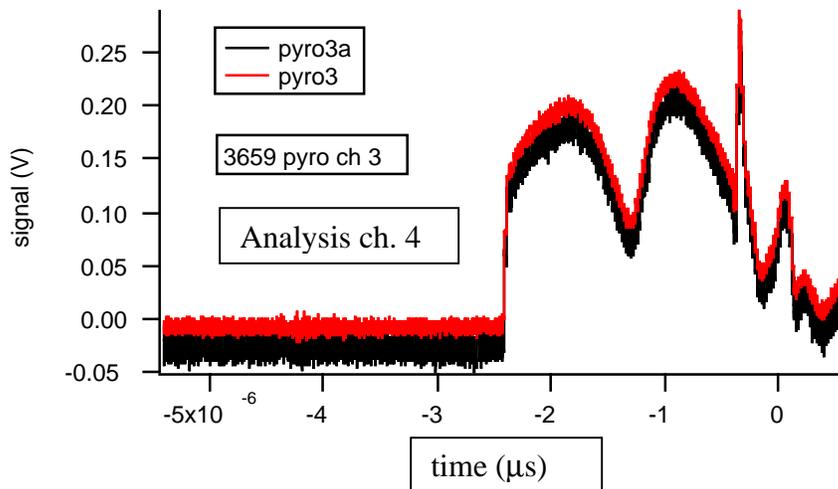


Figure 8c. Signal from analysis channel 4.

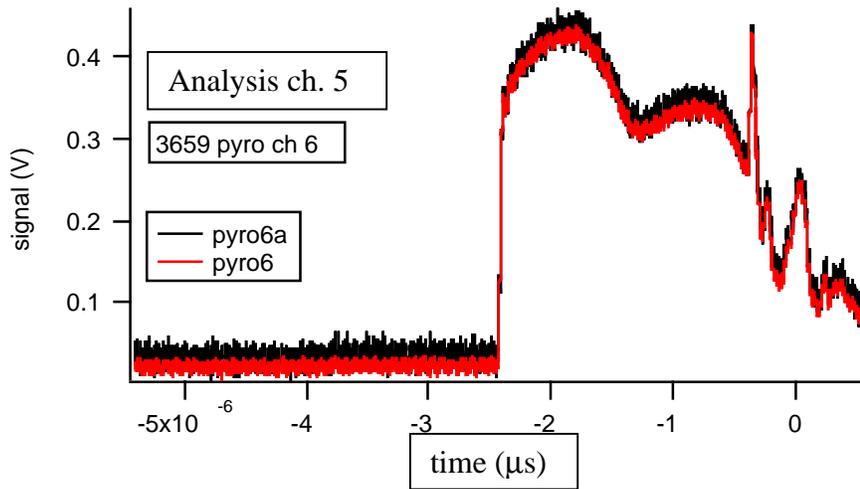


Figure 8d. Signal from analysis channel 5.

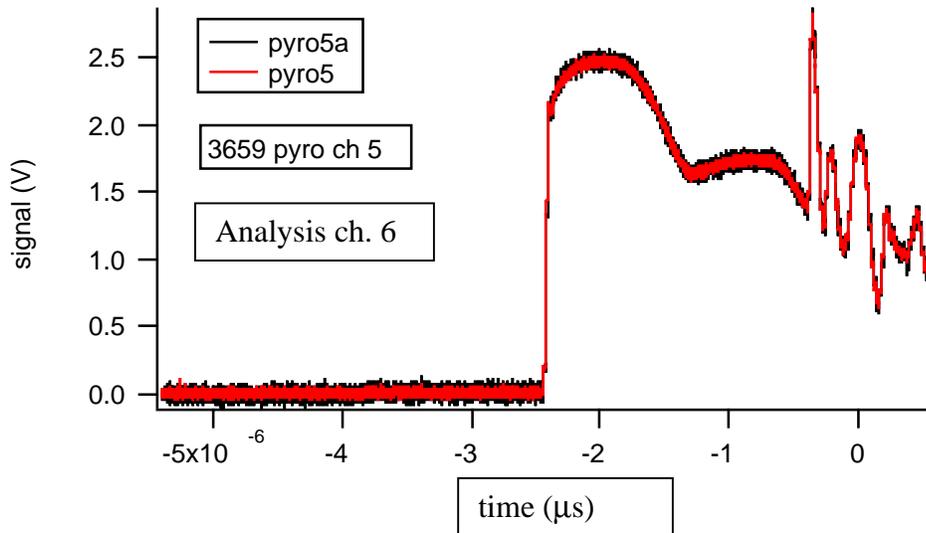


Figure 8e. Signal from analysis channel 6.

There are two difficulties with this data. The first is that the signals are not truly steady state for the expected duration, especially for the longer wavelength channels; the reason is unknown. We have other shots, unfortunately without reflected signals, that do show a steady state. The other difficulty with the shot is that the reflected pulse is late and outside the steady state part of the signal. We won't be able to acquire additional data in the near future, so we will use this shot to illustrate the measurement method.

We take advantage of the observation that analysis channels 5 and 6 have a negligible component of reflected light in the measured signal. We observe that the signal at time -921 ns, which occurs at the peak of the reflected signal, is the same as the signal at time -1400 ns. Furthermore, for these channels, time -921 ns is a time where there is none or negligible contribution of the reflected light to the signal, i. e. the signal is due to emission only. If we now make the hypothesis that the signal due to emission only is the same at times -921 ns and -1400 ns for all channels (observed to be true for channels 5 and 6), we can subtract the signal at time -1400 ns from the signal at time -921 ns and obtain the reflected signal at time -1400. This is obviously not as good as having a shot where it is not necessary to translate the reflected pulse. However, the self-consistency of the computations using this approach supports the hypothesis. We show in table 2 the actual times and the signal levels used in the analysis.

**Table 2. Times and signal levels used in the analysis**

<b>Analysis</b>	<b>time</b>	<b>signal</b>
<b>ch 3</b>	<b>-1.4 us</b>	<b>32.6</b>
	<b>-921ns *</b>	<b>444. mV</b>
	<b>-2.0 us</b>	<b>48.3 mV</b>
<b>ch 4</b>	<b>-1.4 us</b>	<b>118. mV</b>
	<b>-921. ns</b>	<b>232. mV</b>
	<b>-2.0 us</b>	<b>198. mV</b>
<b>ch 2</b>	<b>-1.4 us</b>	<b>28.2 mV</b>
	<b>-921. ns</b>	<b>892. mV</b>
	<b>-2.0 us</b>	<b>23.9 mV</b>
<b>ch 6</b>	<b>-1.4 us</b>	<b>1765. mV</b>
	<b>-921. ns</b>	<b>1729. mV</b>
	<b>-2.0 us</b>	<b>2460. mV</b>
<b>ch 5</b>	<b>-1.4 us</b>	<b>306. mV</b>
	<b>-921. ns</b>	<b>310. mV</b>
	<b>-2.0 us</b>	<b>393. mV</b>

**\*Note: -921 ns is peak of signal in channel 2**

## Detector calibration

The calibration constants  $C_i/C_n$  are obtained by measurements of the radiation emitted by an absolutely calibrated black body source. In the method of analysis advocated here we use calibration factors that are normalized to one of the channels rather than absolutely calibrating each of the channels. We have chosen to normalize to analysis channel 4. The calibration factors are shown in table 3 for each calibration temperature. The factors should of course be independent of temperature and the variation represents an inaccuracy in the calibration. We have used an average for each channel. The variation is not a significant source of error.

**Table 3. Relative calibration factors ( $C_i / C_4$ ), shot 3659**

Channel no.	ch 2	ch 3	ch 4	ch 5	ch 6
<b>Center wavelength</b>	<b>1.239</b>	<b>1.6435</b>	<b>2.253</b>	<b>2.9965</b>	<b>3.982</b>
<b>temperature</b>					
<b>973</b>	<b>3.275</b>	<b>2.336</b>	<b>1.000</b>	<b>0.334</b>	<b>2.128</b>
<b>1073</b>	<b>3.314</b>	<b>2.320</b>	<b>1.000</b>	<b>0.326</b>	<b>2.095</b>

## The Data analysis

We have now shown the elements of the experiment and shown the raw data. The data was analyzed using the analysis and the calibration detailed above. We will first show the results using Method I. In this method we solve for the emissivity in channel 4 and the surface temperature. If the intersections of the curves are not at exactly the same point, an averaging method is used. In this particular example we only use three channels (analysis channels 2, 3, and 4) since the light reflected from the surface in the other channels is insignificant in relation to the emitted light and cannot be determined with sufficient accuracy. Due to the normalization against channel 4, the three channels only produces two curves of the emissivity in channel 4 versus the surface temperature, so there will be only one point of intersection. The curves are shown in figure 9. The point of intersection is at a temperature of 1140 degrees Kelvin and a channel 4 emissivity of 0.20.

Using the temperature of 1140 degrees Kelvin and the channel 4 emissivity of 0.20 we can now use the equations shown previously to solve for the emissivities in channels 2 and 3. They are 0.16 and 0.10, respectively. The complete solution is shown in table 4. It is noted that the emissivity is a strong and irregular function of wavelength. With such a large variation in emissivity the assumption of a gray body emission (emissivities independent of wavelength) is not warranted and would produce inaccurate results. The data will be analyzed below using a gray body assumption to contrast the results.

Table 4. Temperature and emissivities for shot 3659

<b>Temperature K</b>	<b>1140</b>
<b>Emissivity ch. 2</b>	<b>0.16</b>
<b>Emissivity ch. 3</b>	<b>0.10</b>
<b>Emissivity ch. 4</b>	<b>0.20</b>

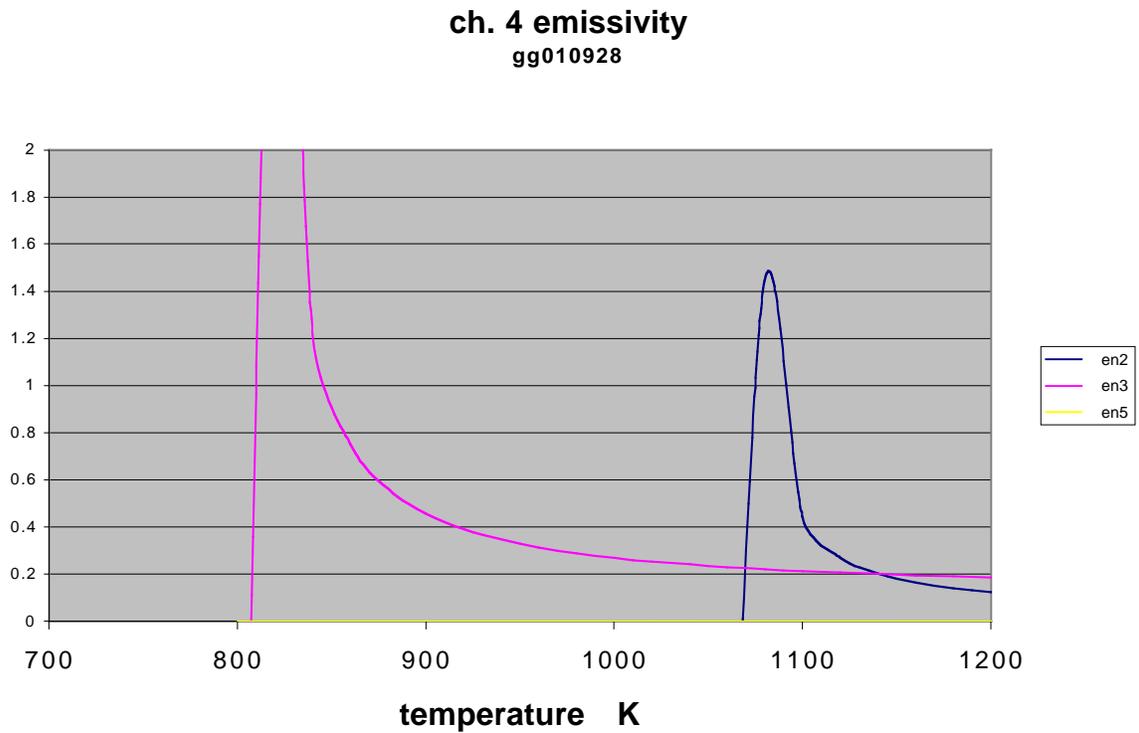


Figure 9. Curves of the emissivity in channel 4 versus the surface temperature in degrees Kelvin for shot 3659. The solution point is at a temperature of 1140 degrees Kelvin and the emissivity in channel 4 is 0.20. The vertical parts of the curves are a plotting artifact connecting negative and positive infinities.

We now turn to Method II. For this method to work we first substitute a highly reflecting surface for the moly surface. It is assumed here that it is perfectly reflecting; its absolute reflectivity can be incorporated in the calculations if it is significantly different from unity. From a direct measurement of the absolute values of the reflected signal from the substituted and the shocked surface in each channel we calculate the ratio of the two values in each channels and obtain the reflectivity in each channel. Subtracting this value from unity yields the emissivities in each channel. They are the same as in table 4 to two significant figures.

To determine the temperature in method II we calculate the measured power in each channel divided by the computed power as a function of temperature. We avoid the use of geometrical factors by dividing this ratio by the corresponding ratio for one of the channels; we use channel 4 again. The solution point should ideally be found at the point where all curves cross unity. In real life there will be some variation and a minimization procedure can be used, similar to that used for the gray body calculation. The plot for Method II for these data is shown in figure 10. The curves cross unity at a temperature of 1140 degrees Kelvin, so the result obtained by this method is the same as for Method I.

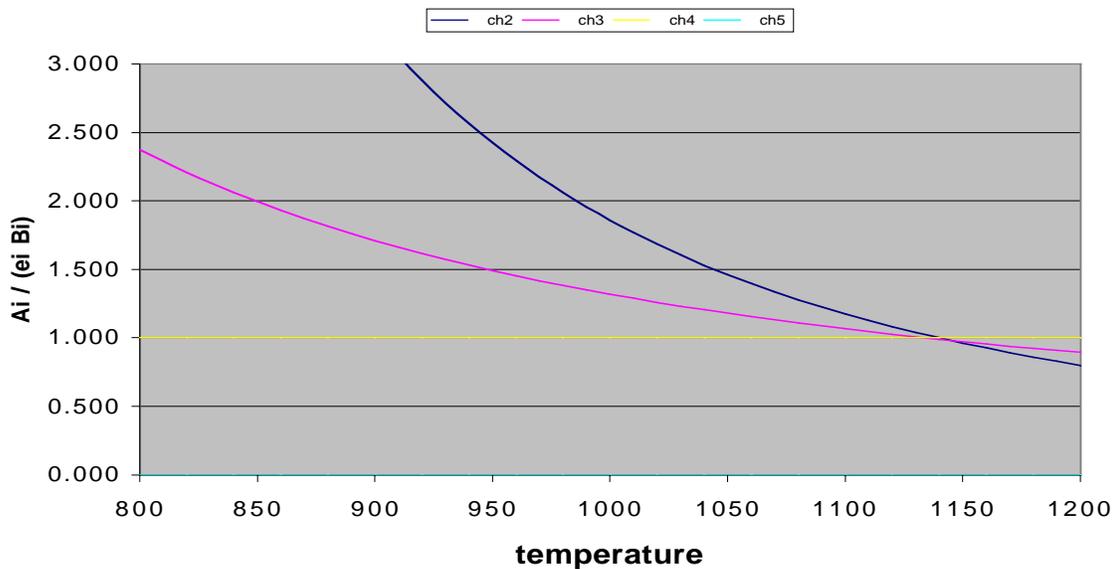


Figure 10. The normalized ratios of measured to calculated emitted power as a function of temperature.

## The Gray body method

An often used method of analysis, the Gray body method, assumes that the emissivity is invariant with wavelength. Only the emission data is used in this analysis. We make a plot of the normalized ratios of measured to calculated emitted power as a function of temperature in a manner similar to that used for figure 10. If the plot shows that the curves meet at the same point, the assumption of equal emissivities is validated. If not, there is no way to know what the temperature really is. Figure 11 shows the gray body plot for the same three channels at the same time used in the analysis using the reflected signals. The figure is directly comparable to figure 10 of method II. It is seen that the intersection denoting a solution point occur at approximately 850 K, 1100 K, and 1450 K. Therefore, the solution points do not validate the assumption of equal emissivities, and the temperature of the surface cannot be accurately determined from these data using a gray body assumption. Using more channels does not aid in finding a solution. Figure 12 shows the Gray body plot for all the channels at time  $-1.4$  us. The spread in temperatures corresponding to intersection points is even greater than in figure 11. We conclude that for these data the Gray body analysis is not valid.

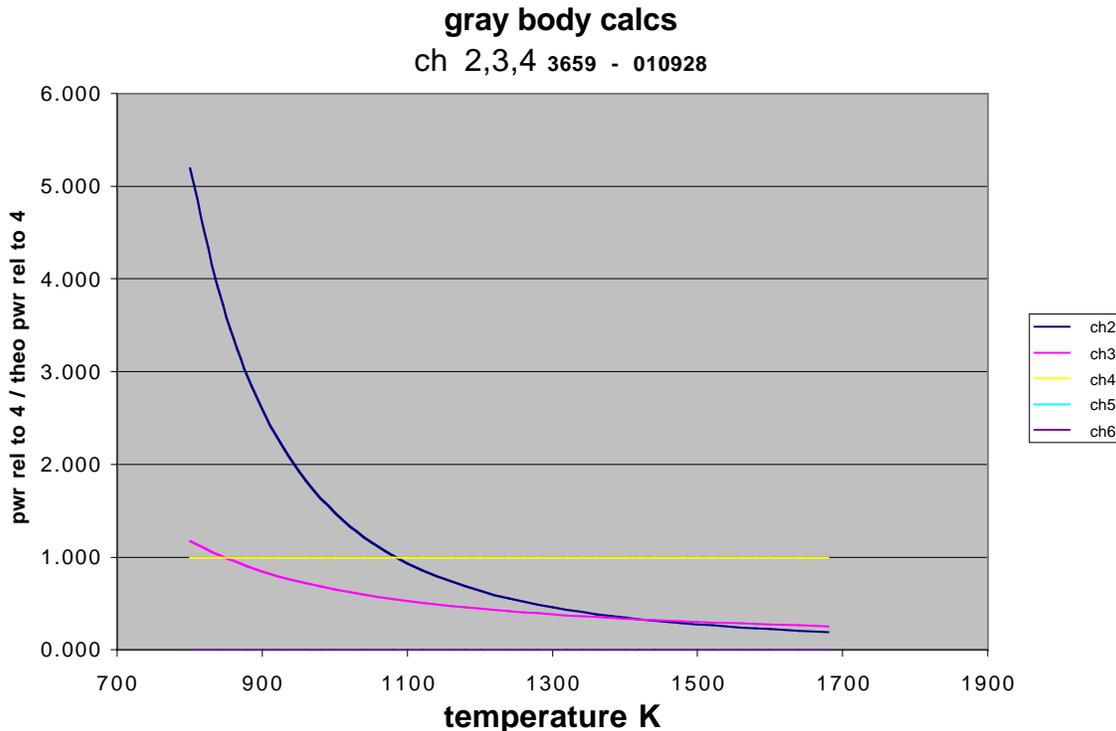


figure 11. Gray body calculations; each of the three intersection is a solution point

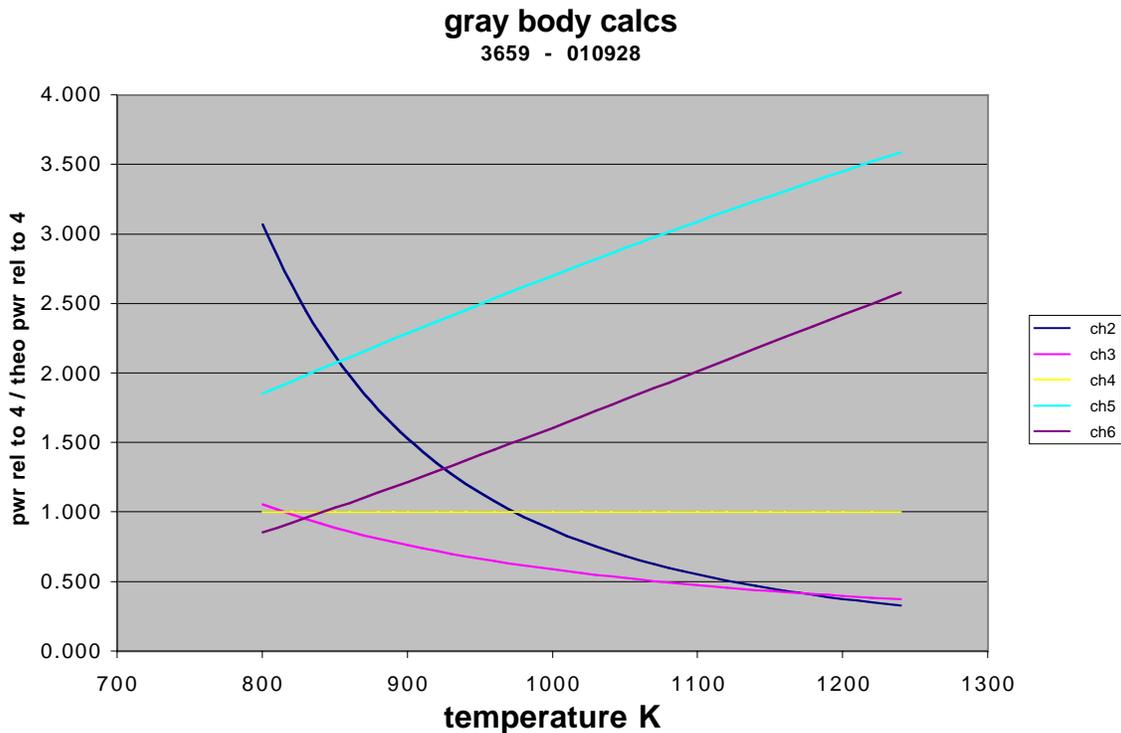


Figure 12. Gray body calculations for all channels

### Future work

A number of improvements can be made to the experiment. The risetime of the incident pulse of light can be made much less than the present  $0.5 \mu\text{s}$ ; times on the order of  $0.1 \mu\text{s}$  should be achievable. The source used for this experiment can be made faster and we have also acquired an alternative source. In addition, it is highly desirable to extend the range of the incident spectrum to greater wavelengths, at least up to  $5 \mu\text{m}$ . The LiF crystal does not pass light very well beyond  $5 \mu\text{m}$ , so another crystal is needed to extend the range. Under some circumstances, the experiment can be done on a shocked surface without a crystal. The reasons for using a crystal for shocked surface experiments is to suppress ejecta and to maintain a high pressure at the shocked surface.

Improvements in risetime and accuracy of timing will diminish the length of time required to have the surface in a steady state. Decreasing the length of time below the present one microsecond will enable the use of thinner targets and decrease the effect of the release at the edge of the impact zone, thus improving the uniformity of the region viewed by the detector.

This experiment used a near non-reflecting enclosure and depended on its success on a geometrical relationship between the emitter, the light source, and the viewing optics. We are also investigating the use of a perfectly reflecting enclosure; the purpose of this enclosure is to integrate the emission over angles and to increase the effective emissivity of the surface. The latter feature is useful under conditions where the light emitted is otherwise marginal for measurement.

## **Conclusion**

We have shown that the channel emissivities and the temperature of a shocked surface can be determined accurately and self consistently by measuring both reflected and emitted light in each band of wavelengths of a multichannel spectrometer. At least three bands are required. Two independent methods have been used. One method uses only the knowledge of the spectrum of the incident light as well as the emitted and reflected light. The other method requires a calibration of the reflected light from a surface having a known reflectivity (such as gold) in an identical geometry. The first method can be used when the emissivities vary sufficiently with wavelength. If the emissivities do not vary, the Gray body assumptions can be used for the analysis and will yield the correct results. The second method can be used under any conditions, assuming, of course, that the measurement of the known surface can be accomplished.

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University of California  
Lawrence Livermore National Laboratory  
Technical Information Department  
Livermore, CA 94551

