
Scanning Parameter Space for NIF capsules in HYDRA



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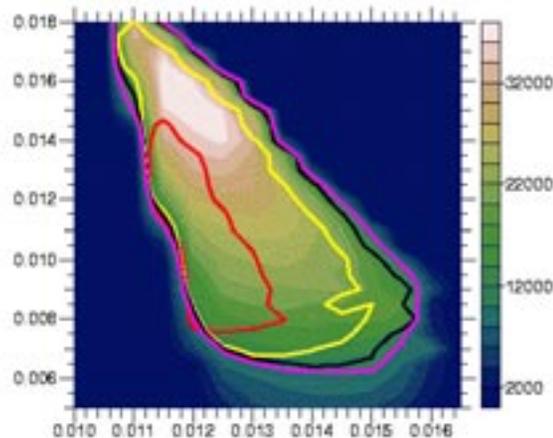
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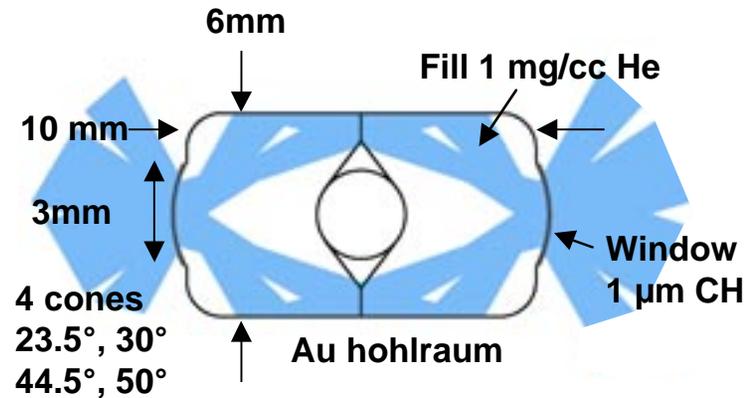
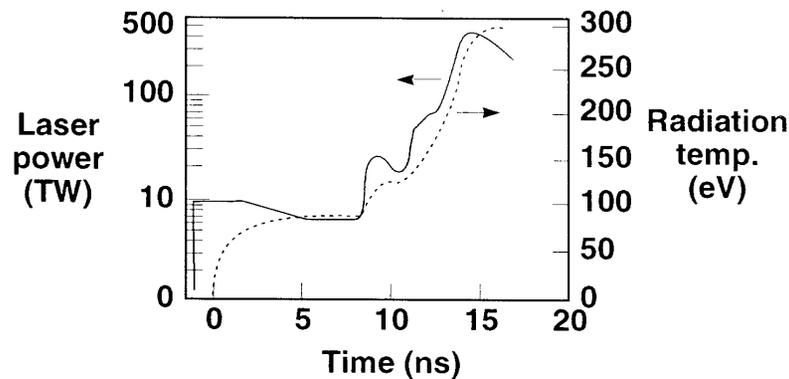
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Summary



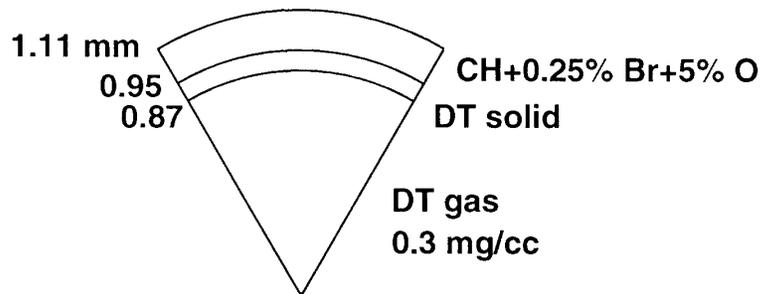
- We have implemented an automated pulse shaper for NIF capsules in HYDRA.
- We have developed the infrastructure to do scans using the automatic pulse shaper across any n-dimensions of capsule parameter space.
- Using this infrastructure, we have performed several scans examining parameters for uniformly doped Beryllium capsules.
- To coordinate more closely with the anticipated experimental shock timing strategy, we have started to develop an automated pulse shaper which uses planar geometry and liquid DD.

We are interested in choosing the optimum capsule for ignition experiments



There are at least eight critical values in pulse shaping, including timing, amplitudes and total flux

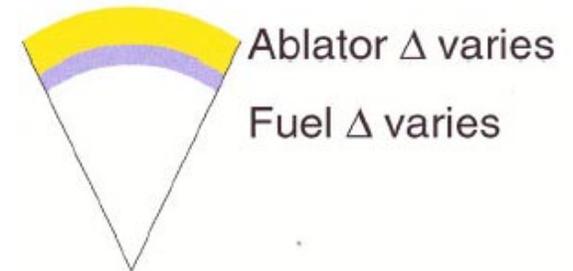
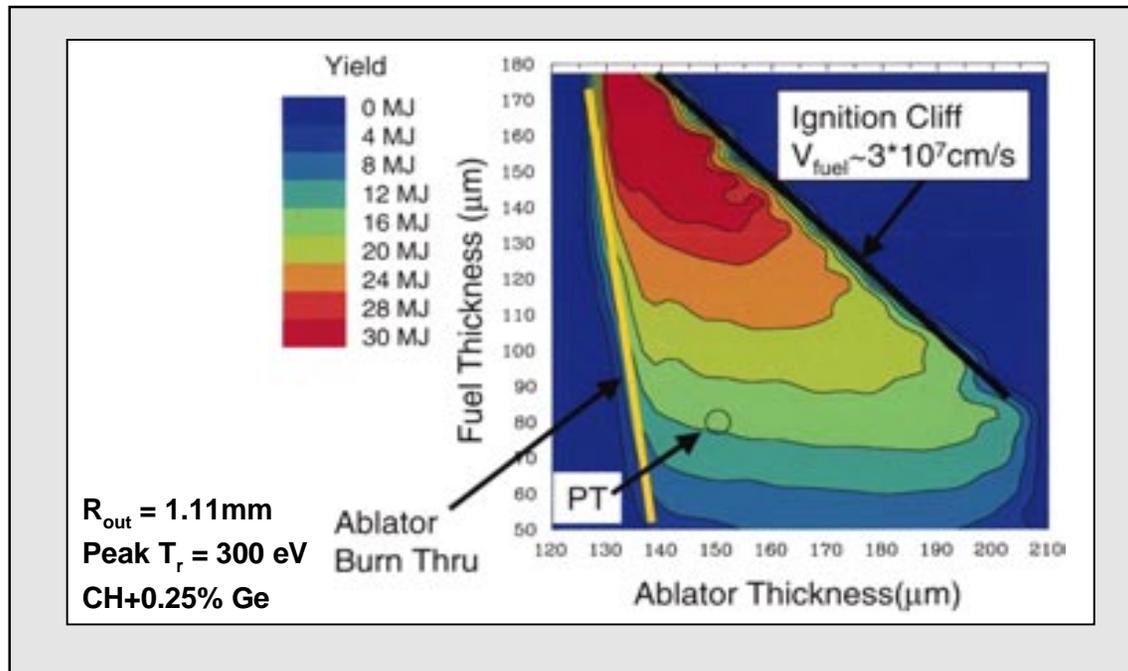
Hohlraum design can affect temperatures, spectrum, and many other variables



Capsules have different sizes, different materials, and different material thicknesses

There are many variables to choose from in capsule design!

Some scans of capsule parameter space have already been done



- This contour map is made from a suite of ~ 500 1D capsule designs (5 μm spacing)
- Maps the space of performance according to 1D simulations

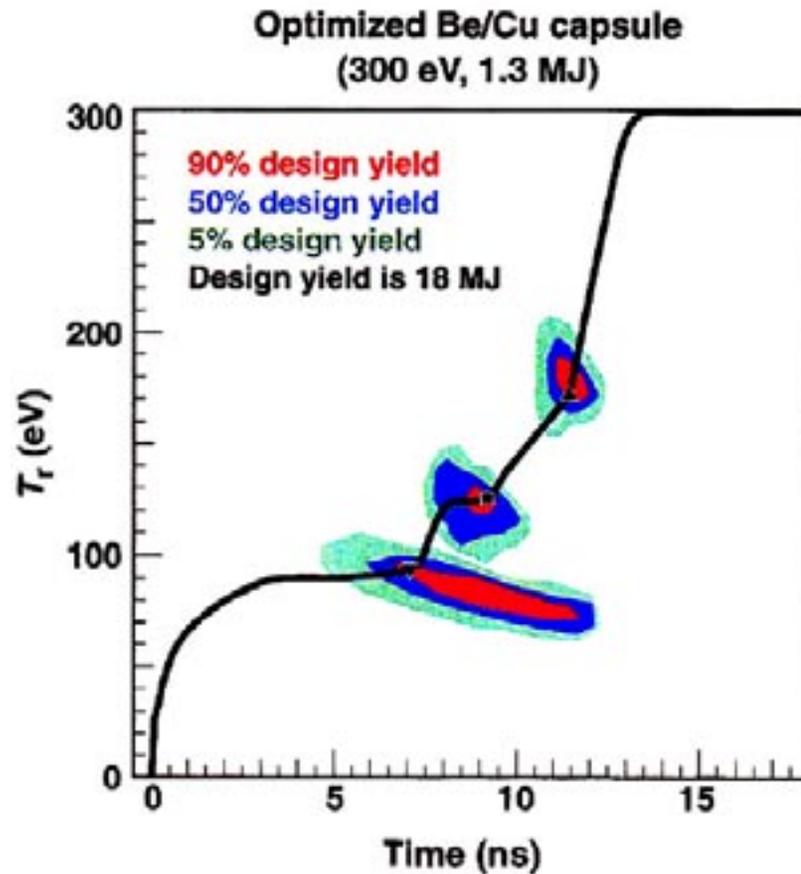
We want to develop a better scanning ability



While the current scan technology has been useful, we would like the capability to:

- **Scan using the code HYDRA**
- **Make changes to any input parameter without requiring significant change to the code**
- **Perform >2 dimensional scans, over parameters other than ablator or fuel thickness**
- **Use a pulse shaping technique more similar to anticipated experimental methods**

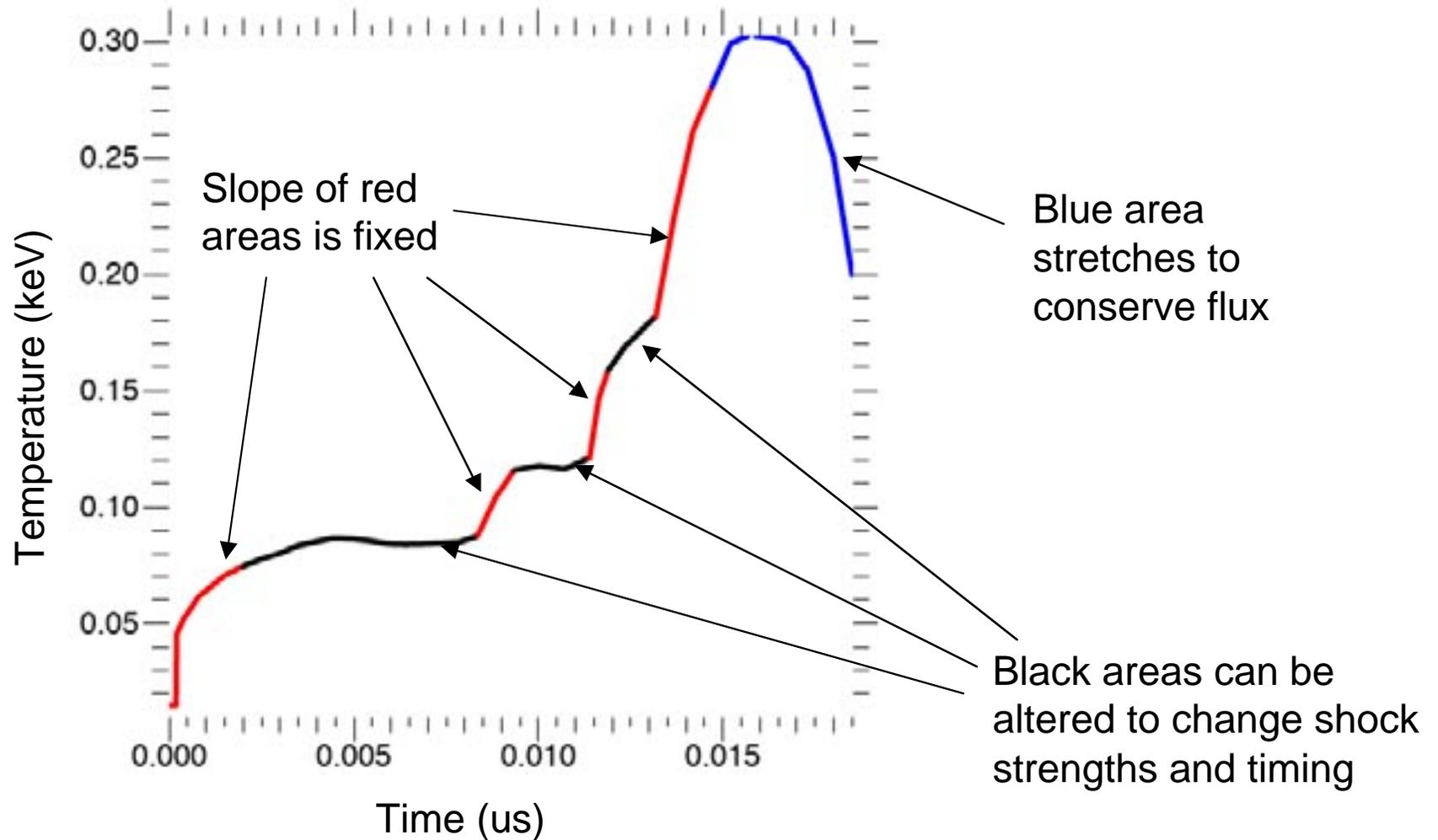
Capsules are very sensitive to the pulse shape



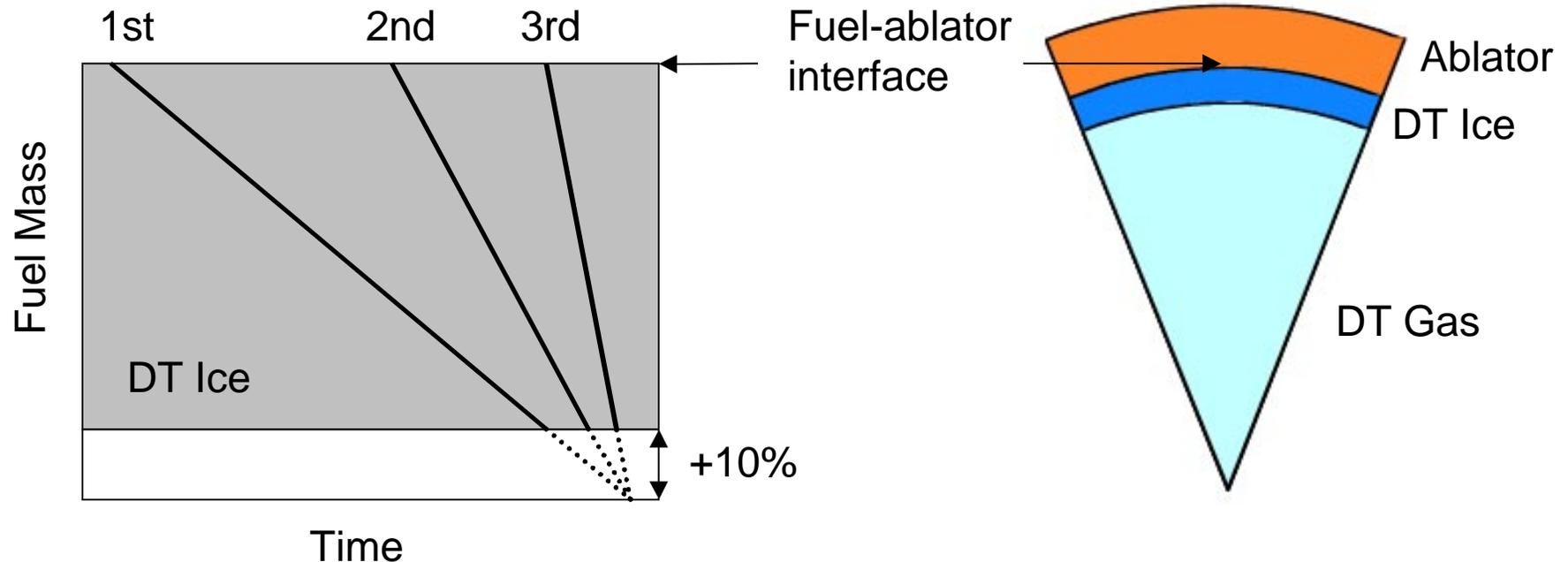
- The “banana plot” to the left shows that single shock shifts by as little as 100 picoseconds can be detrimental to capsule performance
- Shock strengths are not as vulnerable, as long as timing is changed accordingly.
- This diagram was produced by Dave Munro

The first step is to develop an automated pulse shaper

Shock timing can be changed through modifying the radiation drive

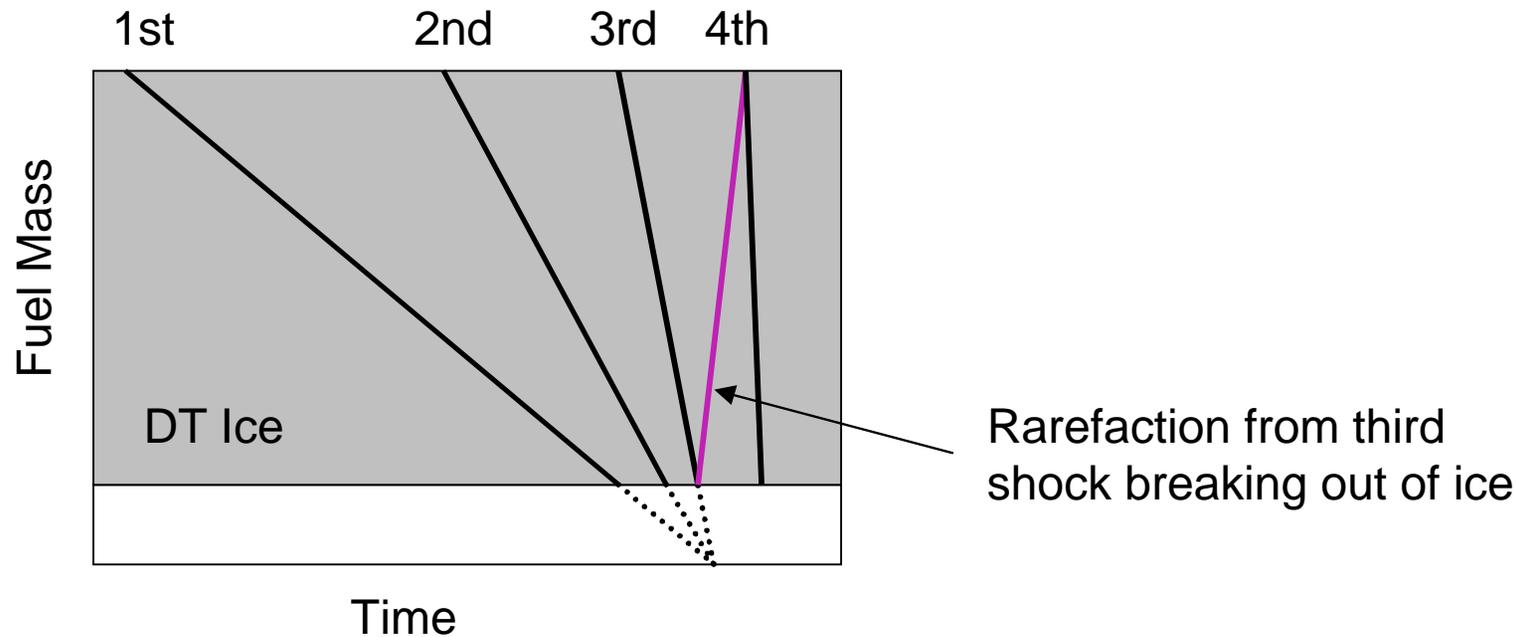


Shock timing is determined by Dave Munro's coalescence criterion



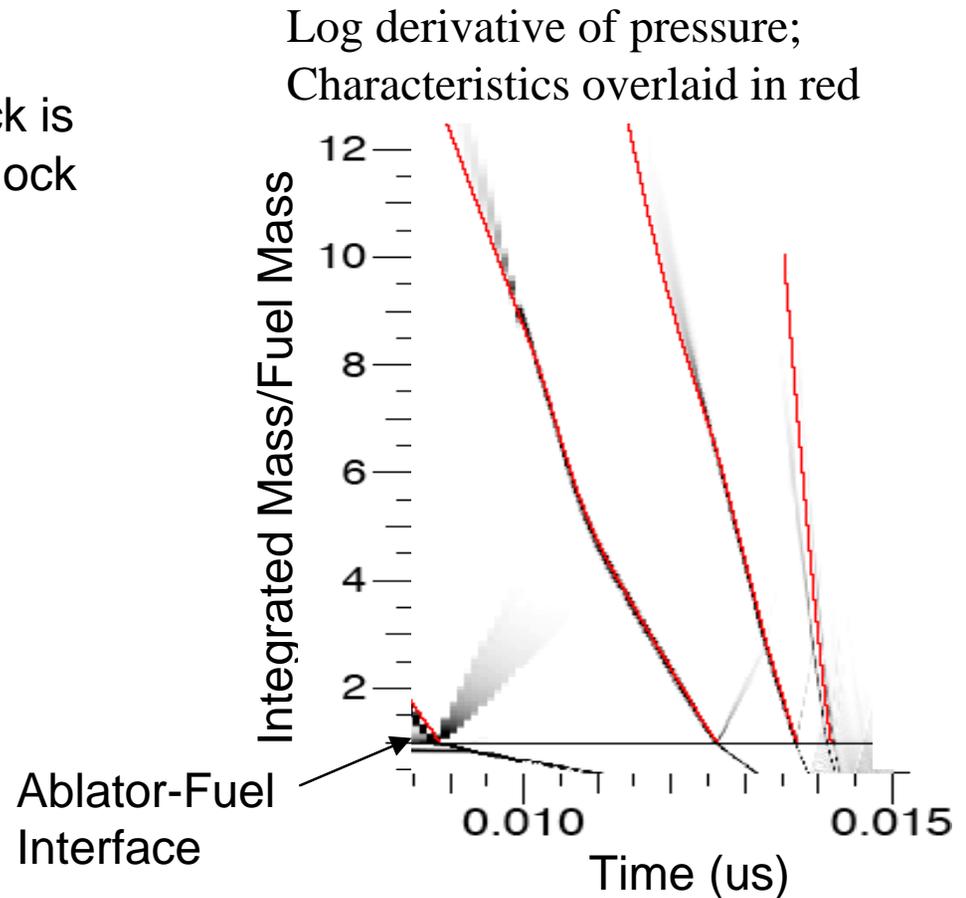
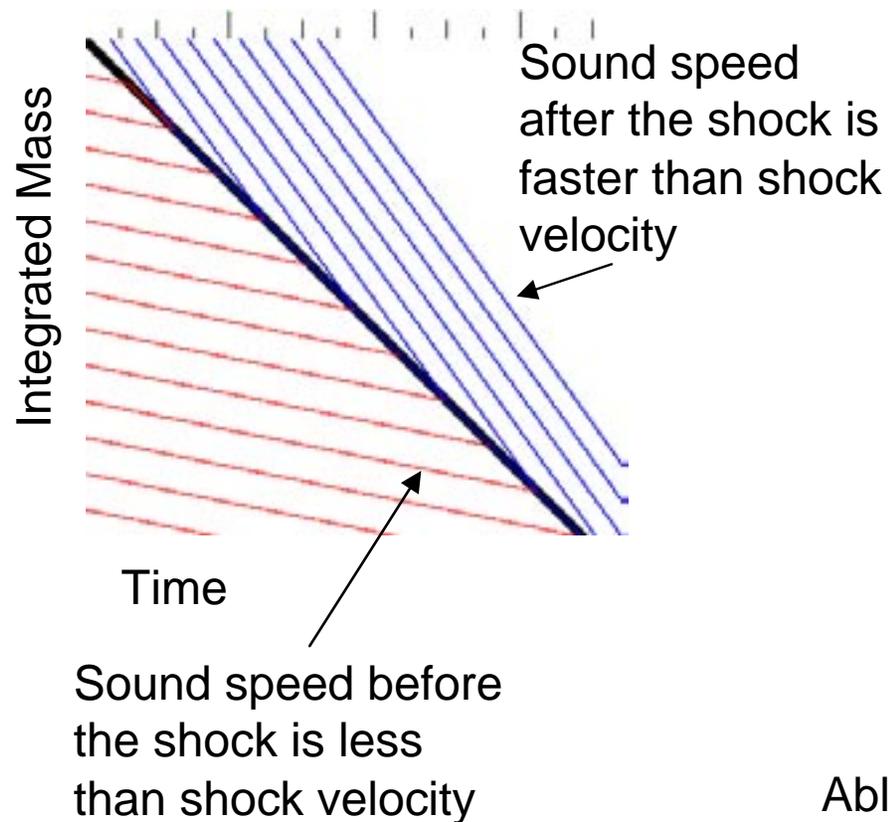
- The goal is to have the shocks “break out” in close succession, but not run into each other too early.
- Shock velocities are determined by pressures specified in advance
- Then, timing is found by working backwards from the coalescence depth, usually at 110% of fuel thickness

Fourth shock timing is done differently



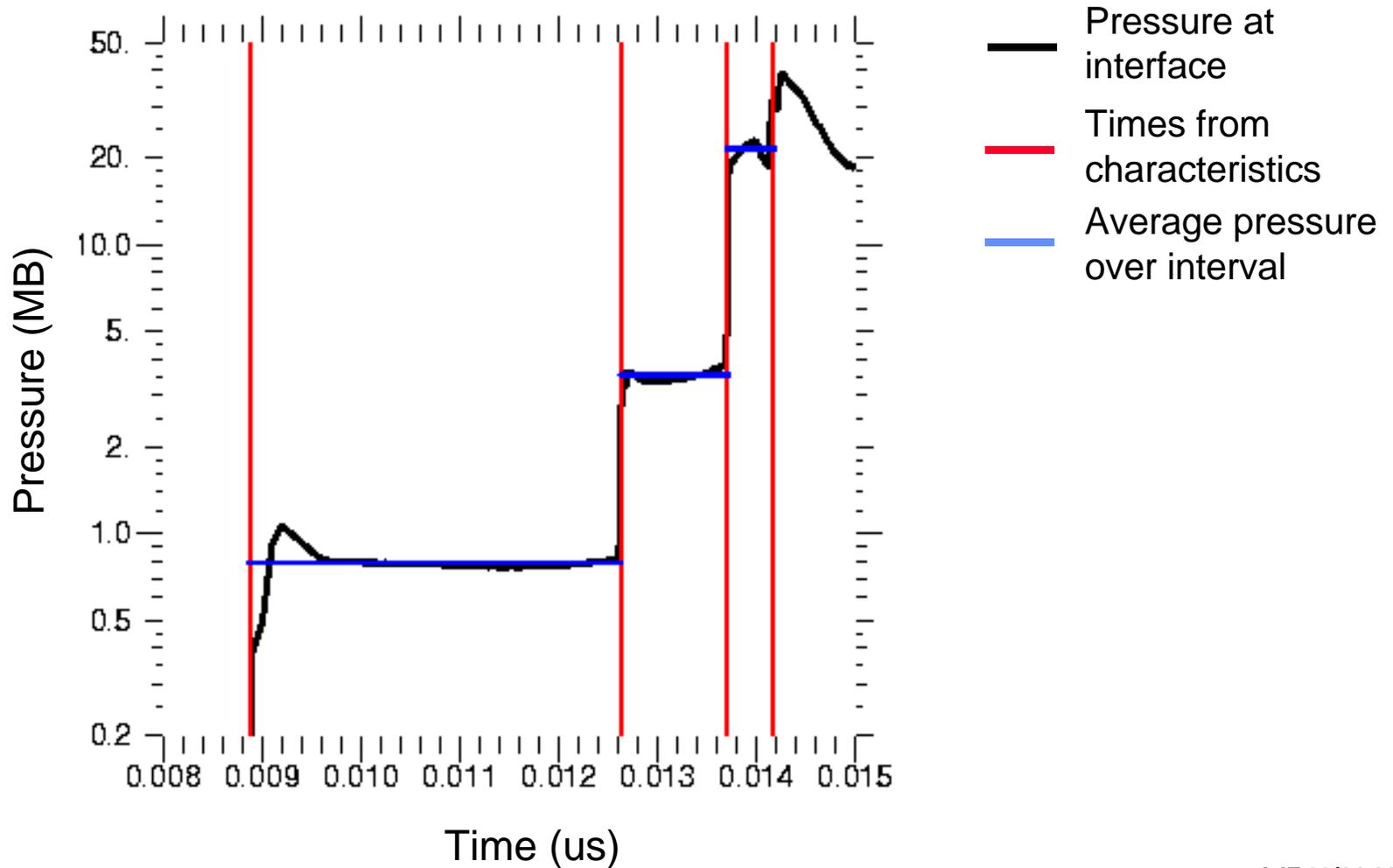
- A rarefaction is launched when the third shock breaks out of the fuel
- We want the fourth shock to cross the fuel interface at the same time as this rarefaction

Characteristics can be used to find shocks



A postprocessor capable of tracking these characteristics was designed and implemented

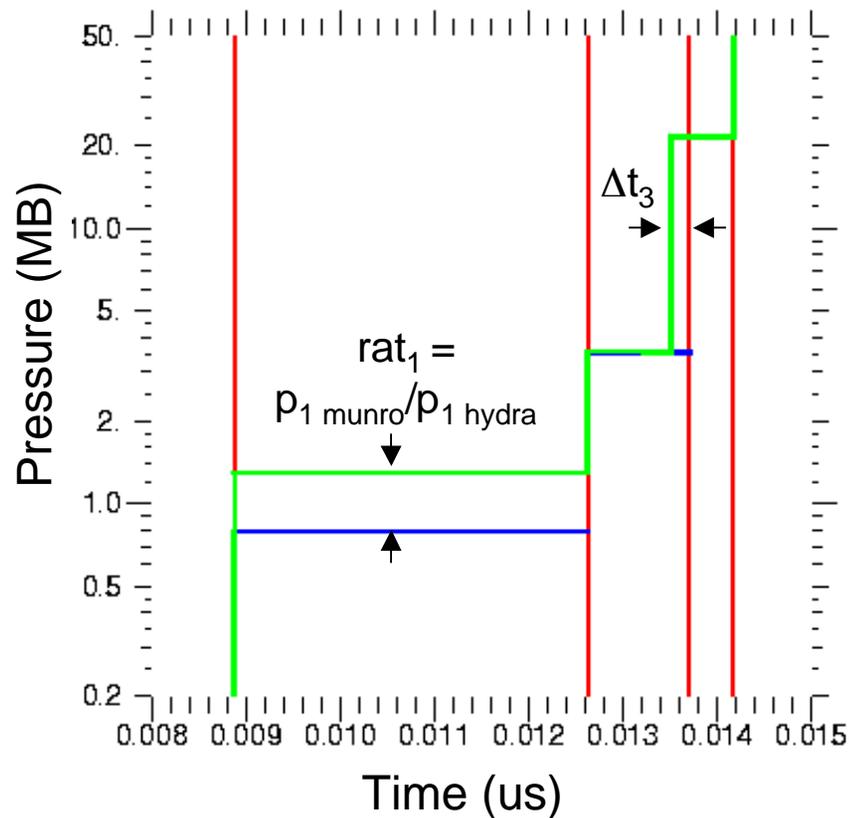
The pressure profile at the fuel-ablator interface is then analyzed



The pressure profile is compared with the Munro criteria and appropriate changes are made

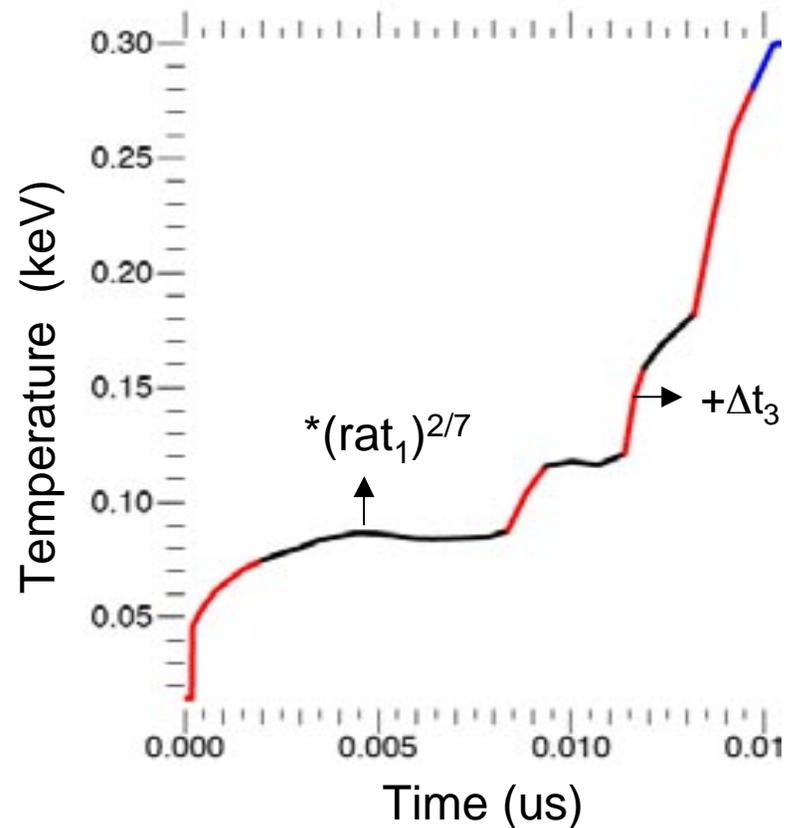


Differences between the target pressure profile and actual profile are analyzed . . .

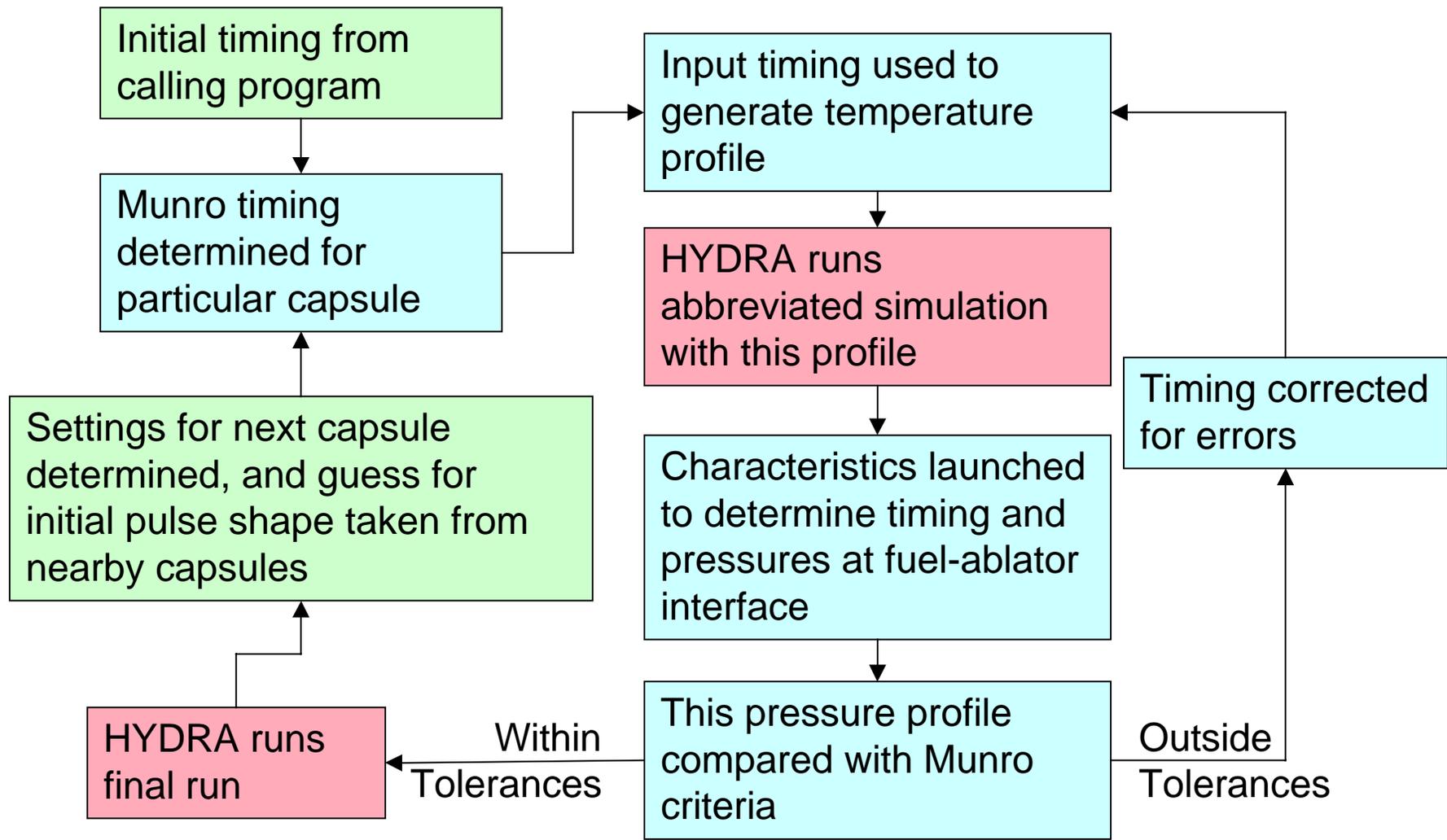


— Target pressures and timings

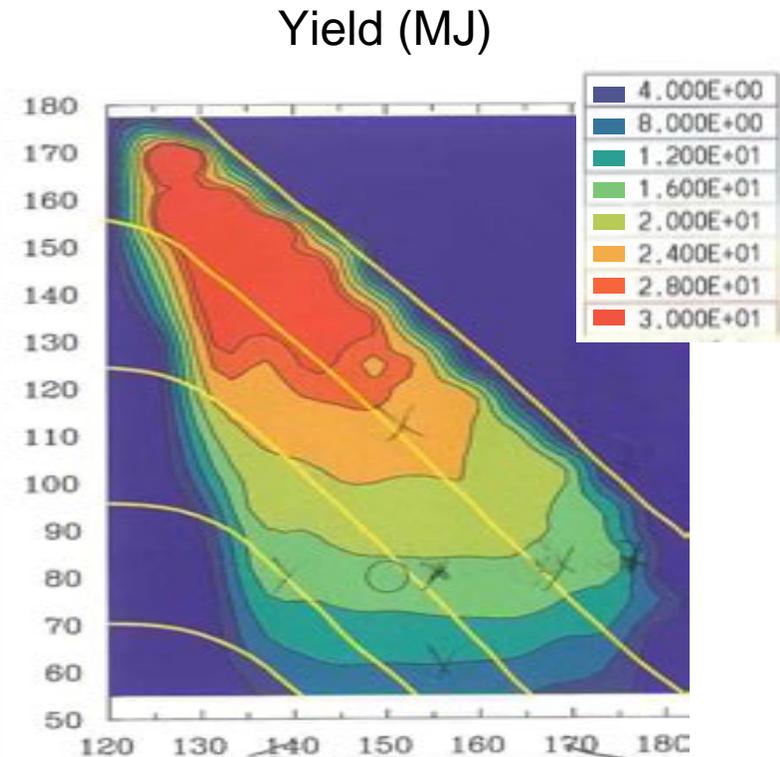
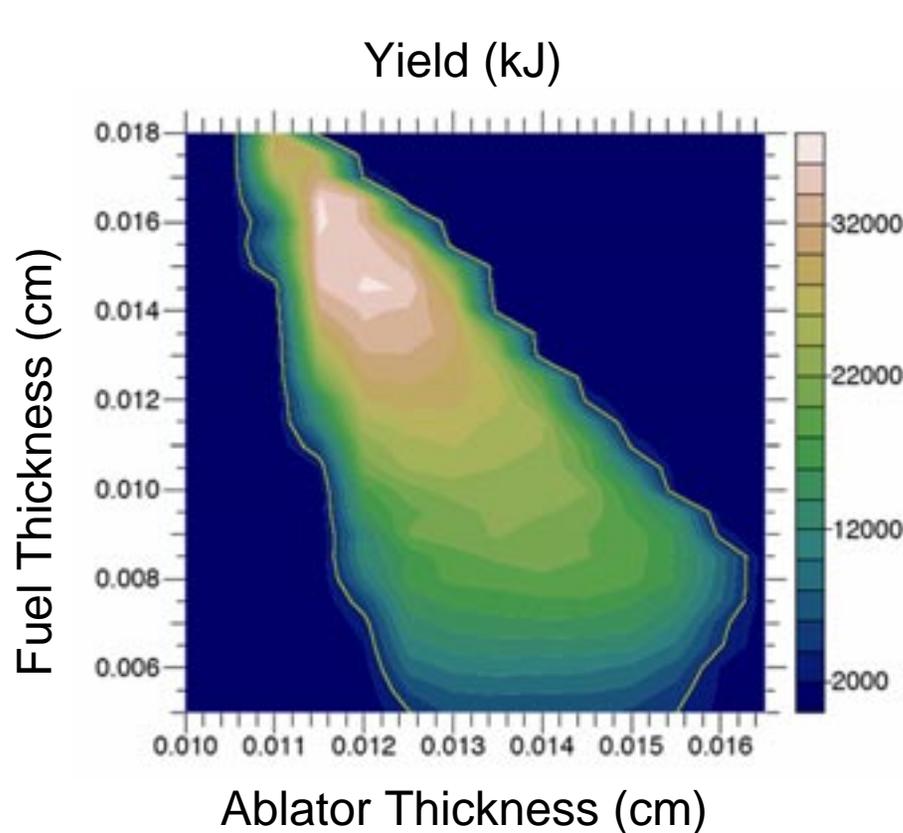
. . . and necessary changes are made to the radiation profile



This process is iterated until a proper pulse shape is found



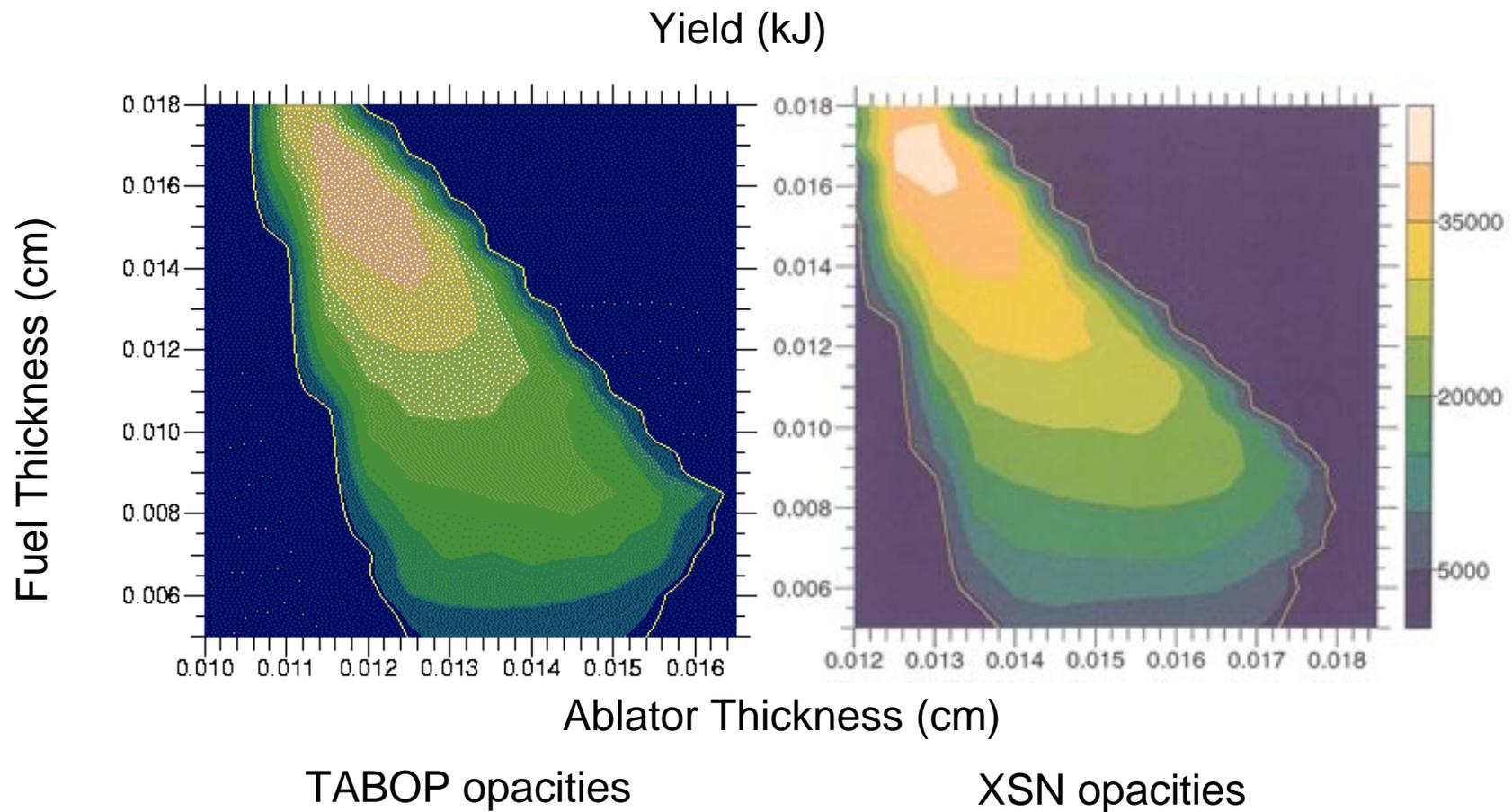
Our results show reasonable solutions for .9% Cu doped Be shell capsules



Mark Herrmann's LASNEX
Results for the same capsule

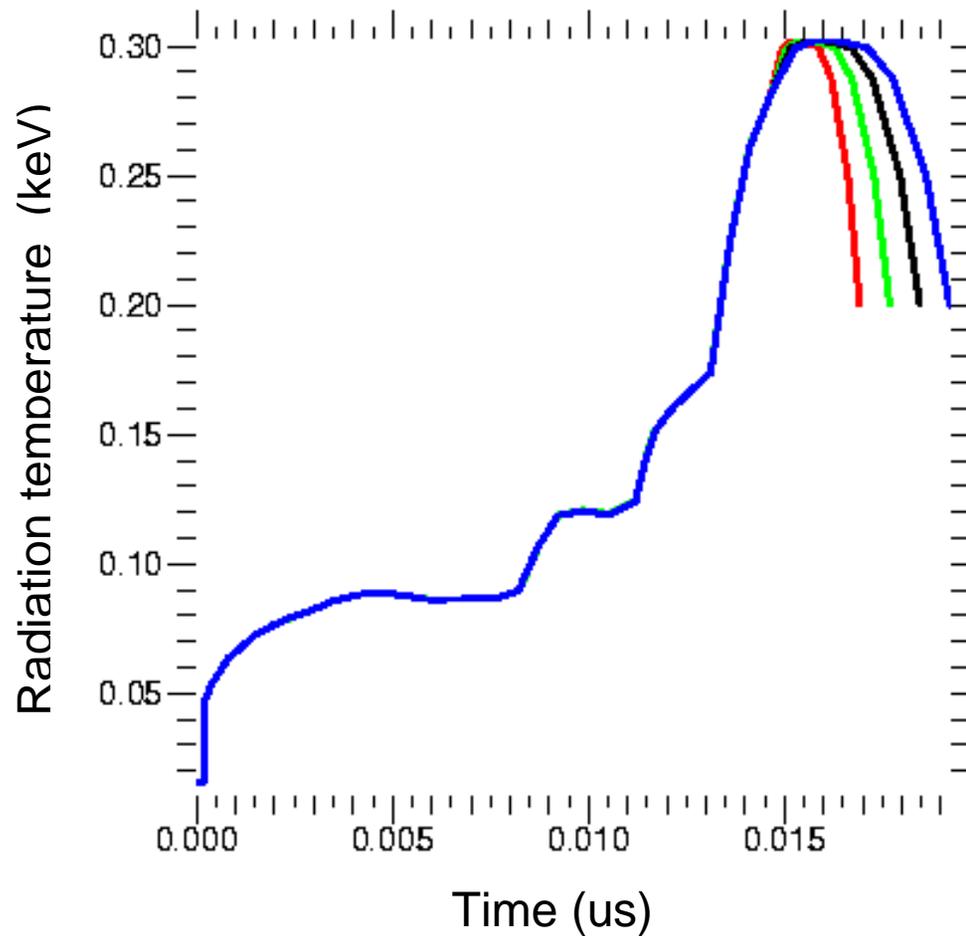
Note the ~20 μm "shift" in Ablator Thickness.

The space shifts depending on opacities



This 20 μm shift is due to differences in opacities

Changes in total fluence can have an important effect on the ability of the capsule to burn

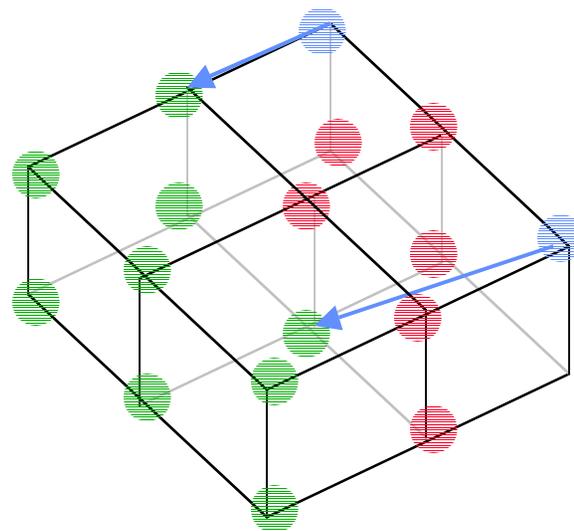


- 60% of Nominal
- 80% of Nominal
- Nominal NIF Drive
- 120% of Nominal

$$\text{Fluence} = \int T_r^4 dt$$

Laser energy \approx Fluence

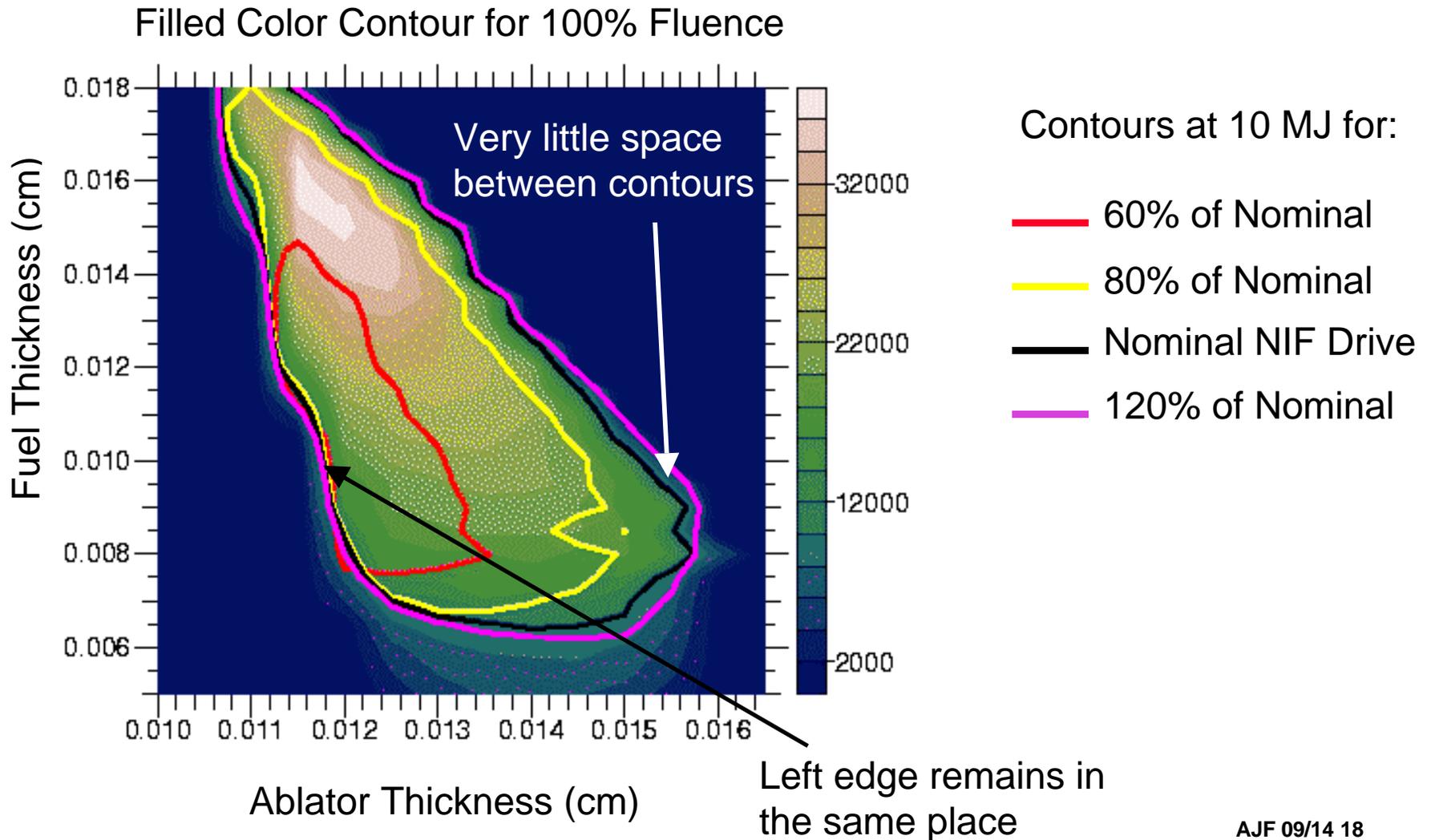
Starting runs use the timing from the nearest finished run in n-dimensional space



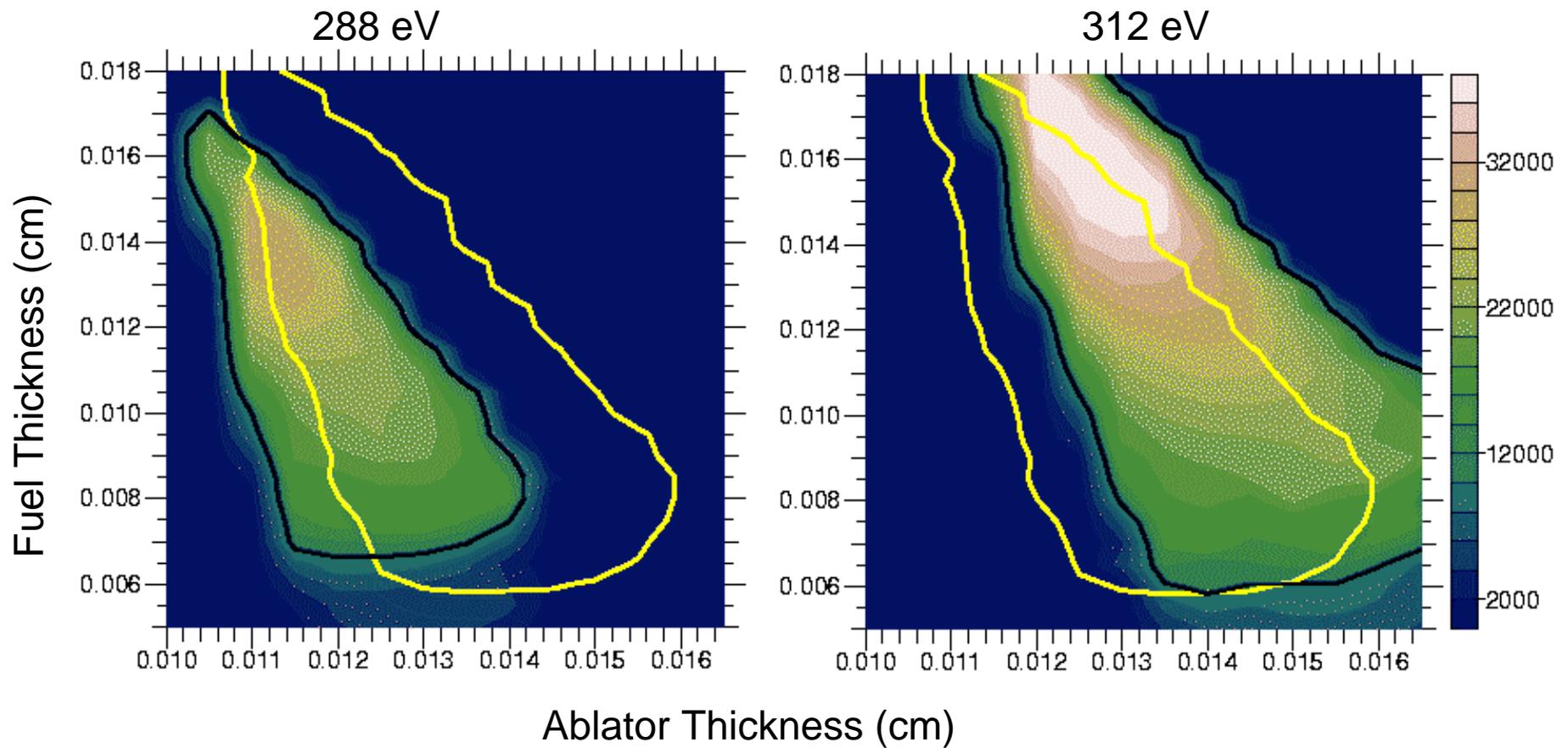
- Finished runs
- Unfinished runs
- Starting runs
- Path to nearest finished run

Scanning over n-dimensions is much more efficient than repeating 2D scans

While raising the fluence produces few benefits, reducing the fluence can impact performance



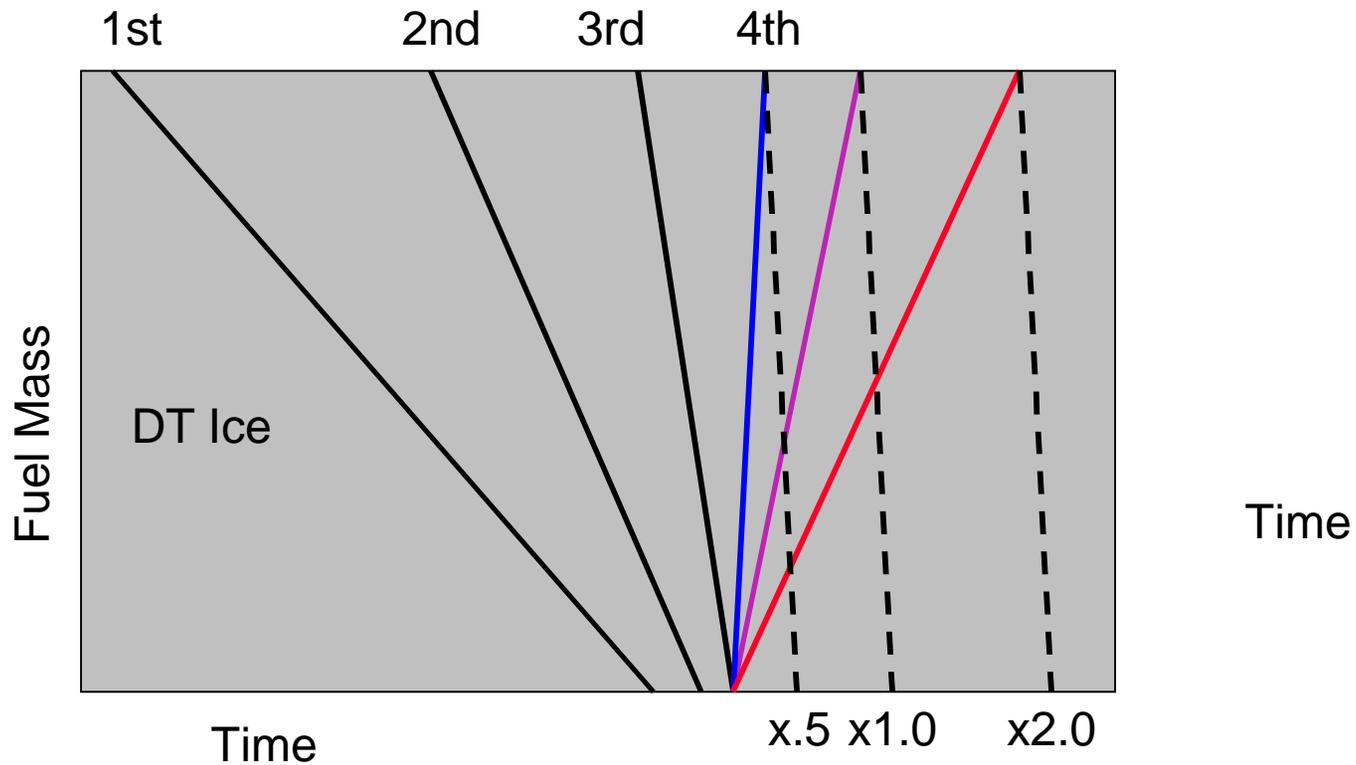
We also found that peak radiation temperature has an impact on the design space



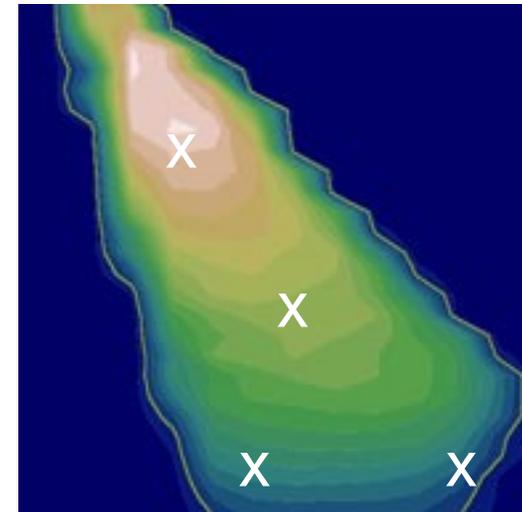
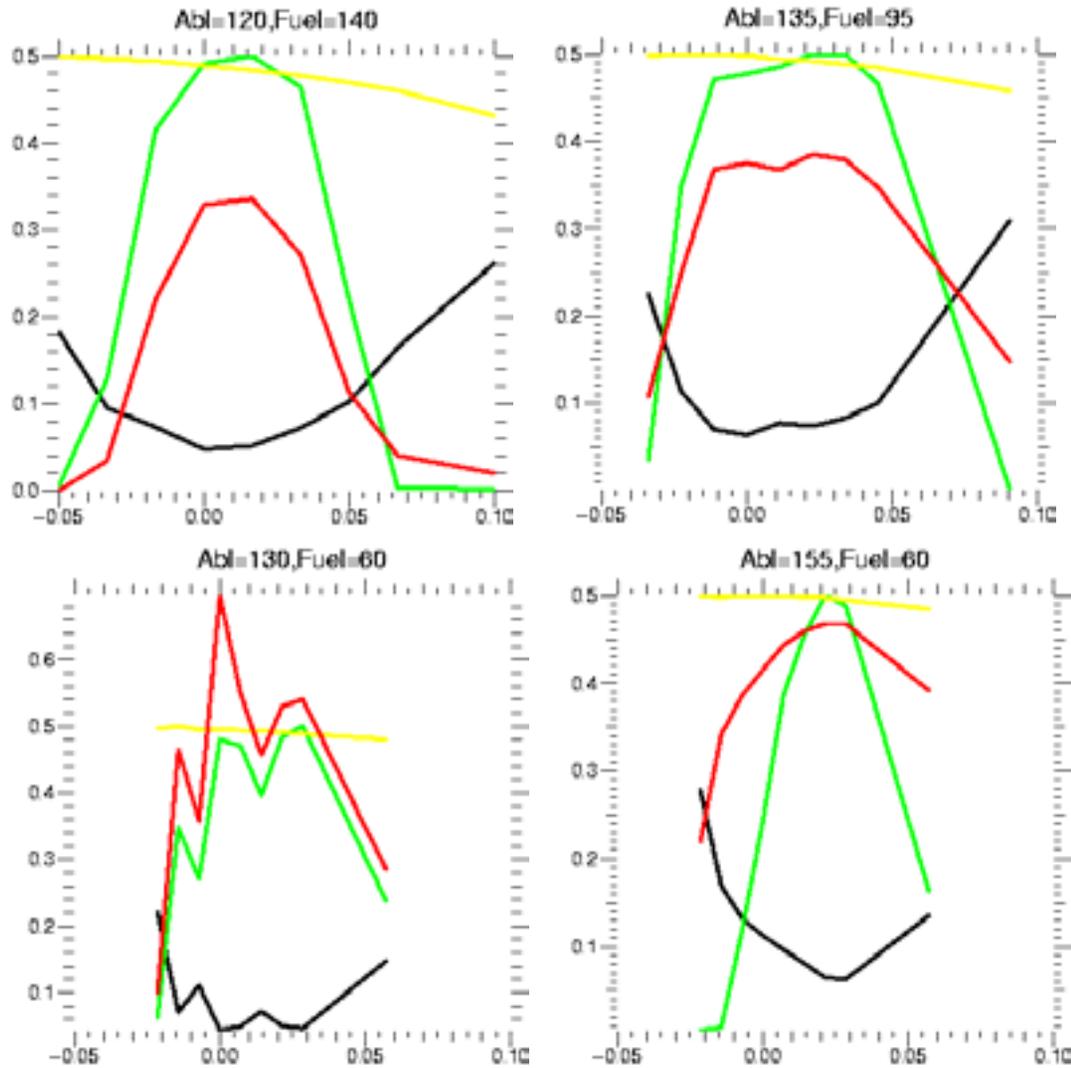
— 10 MJ Contour — 10 MJ Contour for 300 eV

Total flux is the same as the 300 eV capsule

To better understand fourth shock timing, we conducted a scan over the fourth shock space



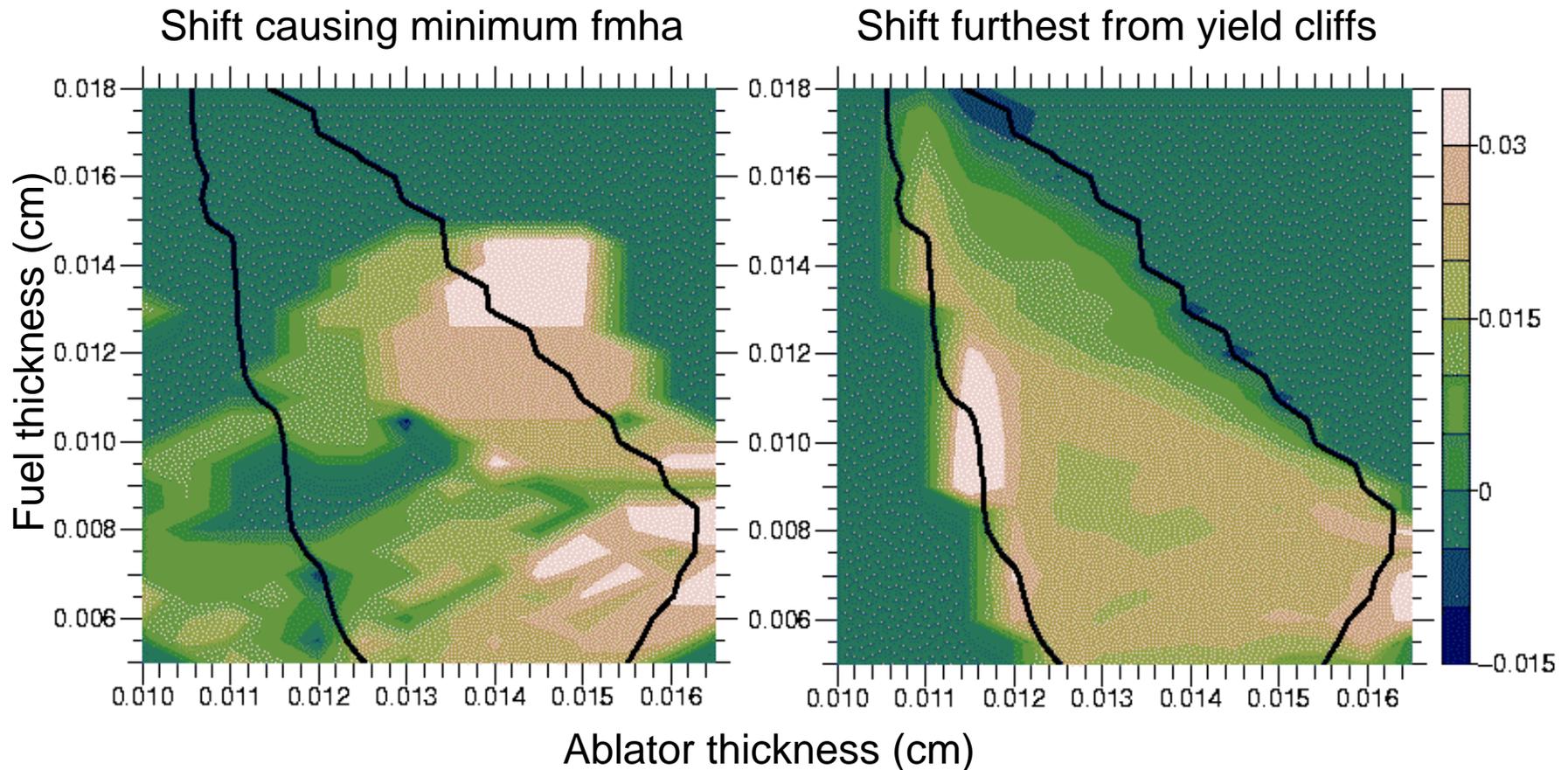
Scanning over fourth shock timing show some improvements can be made



- Normalized velocity
- Normalized yield
- Margin
- FMHA

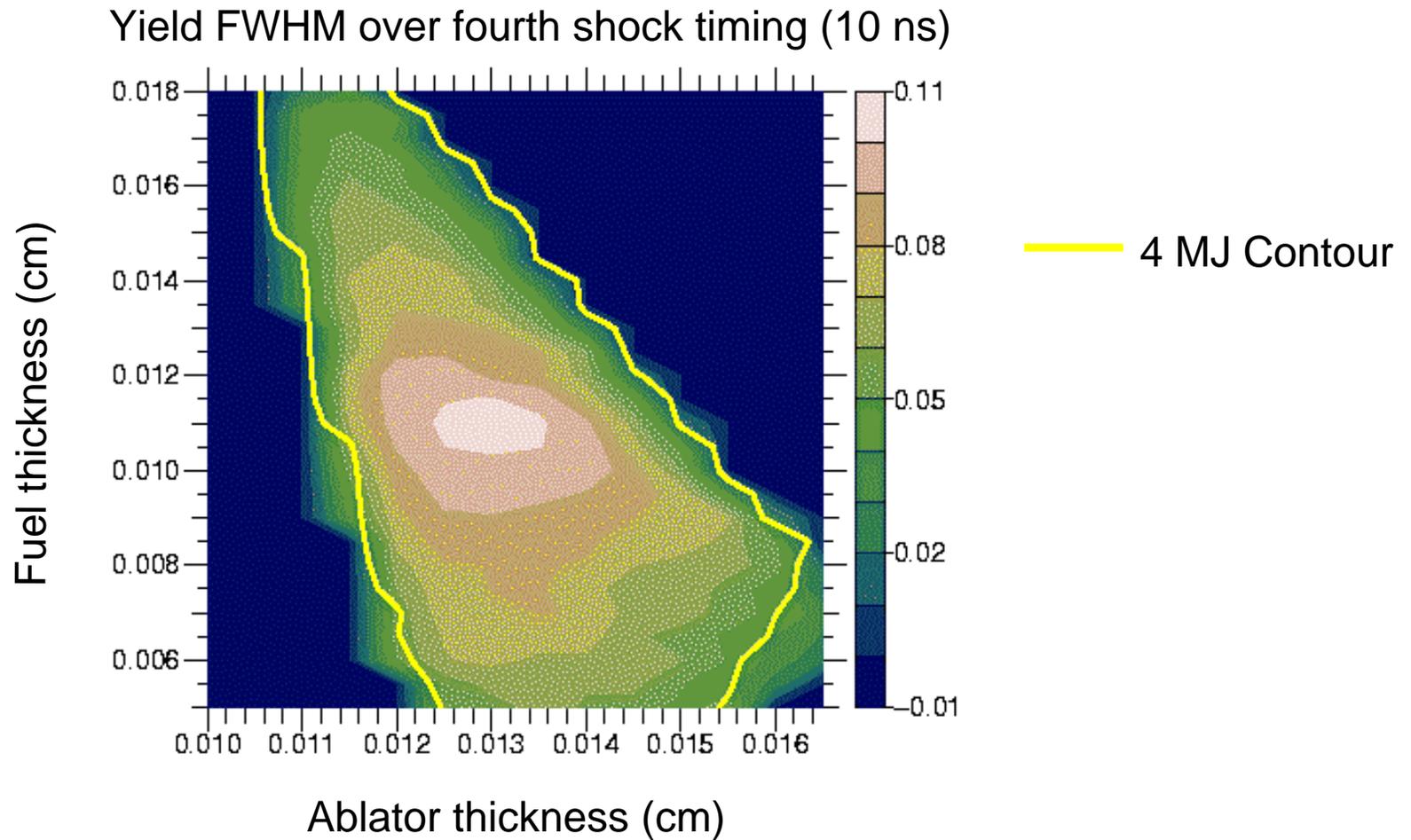
Fourth shock time shift (10 ns)

We can find the optimum fourth shock timing over the interesting space

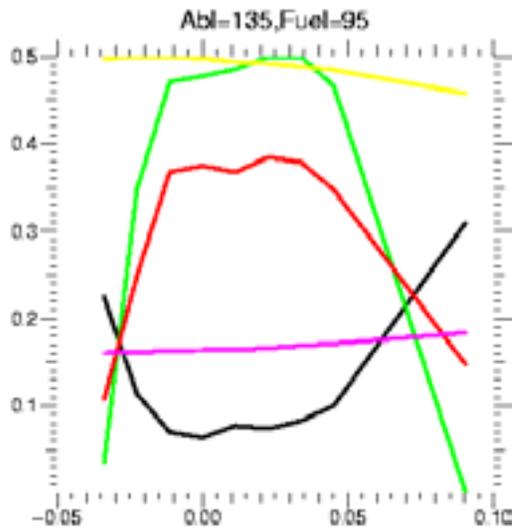


While the methods differ, and vary over the space, an optimum shift for the most interesting area is about 200 ps

Tolerance to changes in the fourth shock timing is best in the center of the space



The fourth shock timing is not the only thing changing when the target time is moved



Fourth shock time shift (10 ns)

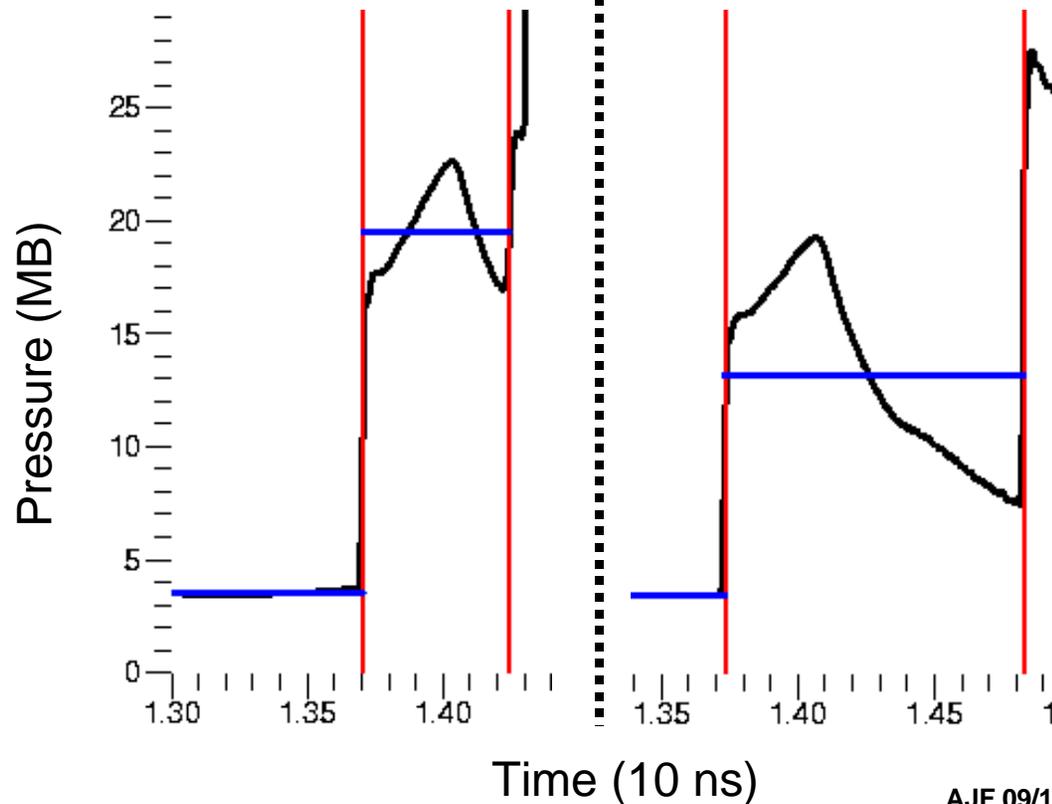
— Third shock launch temperature

In response to this apparent DECREASE in amplitude, the pulse shaper INCREASES the third shock temperature

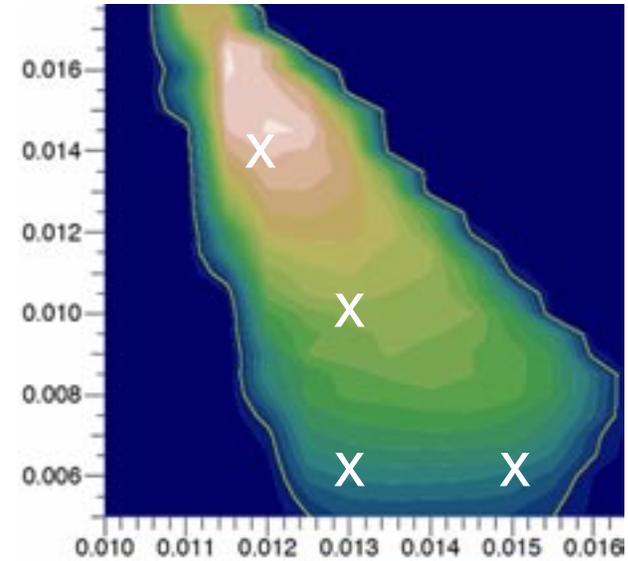
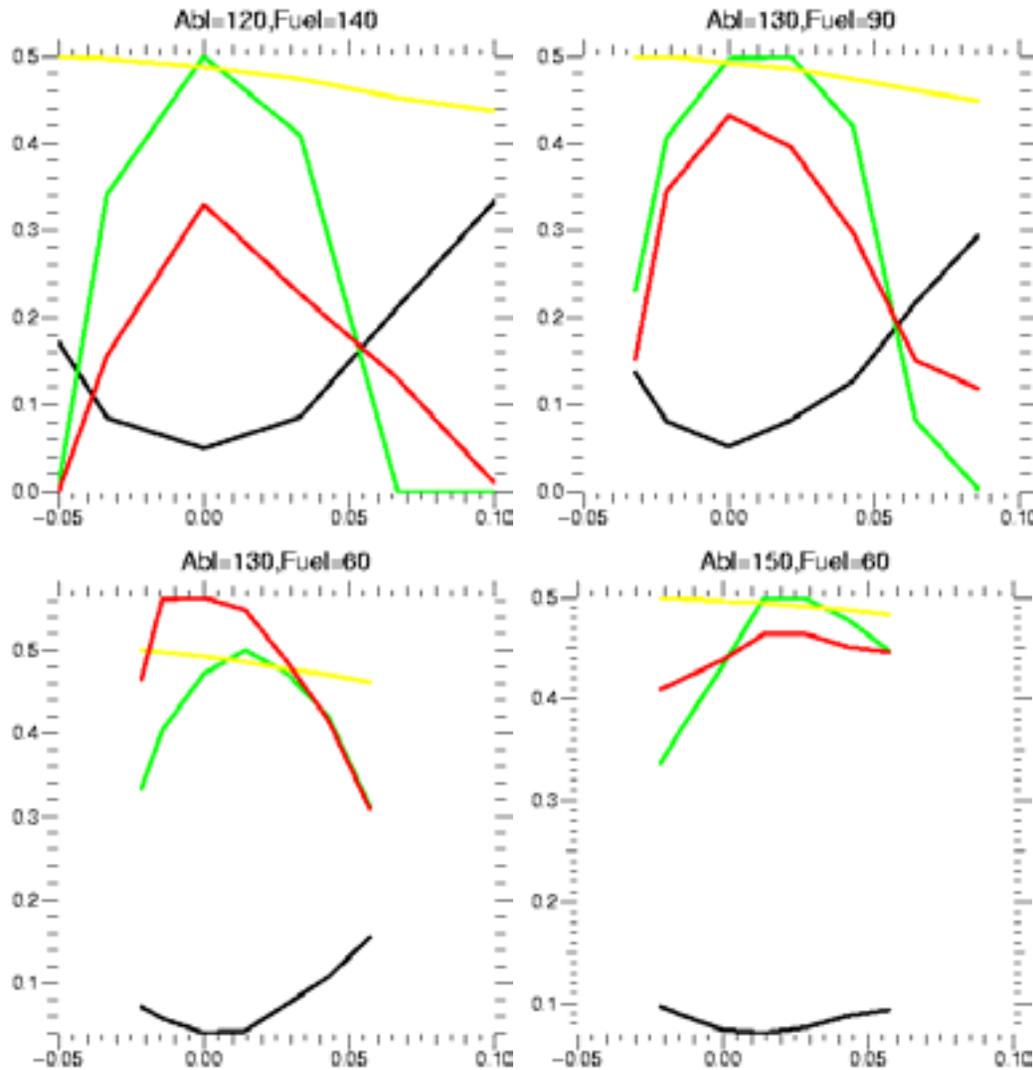
Pressure profile for third shock at fuel-ablator interface

Fourth shock timing:

Normal timing -200 ps Normal timing +200 ps



Shifting the fourth shock without modifying the third yields similar results



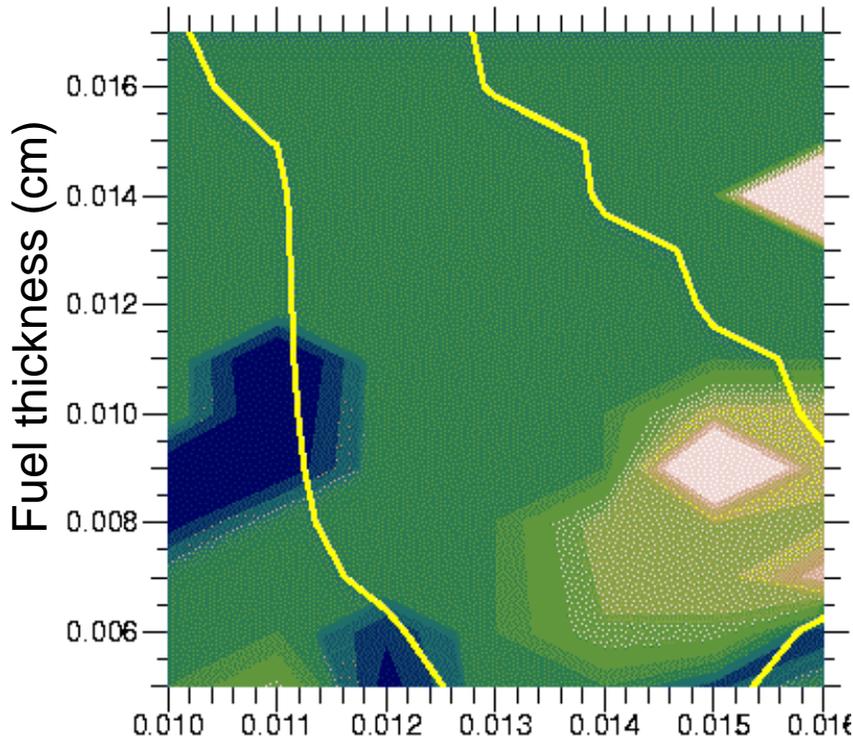
- Normalized velocity
- Normalized yield
- Margin
- FMHA

Fourth shock time shift (10 ns)

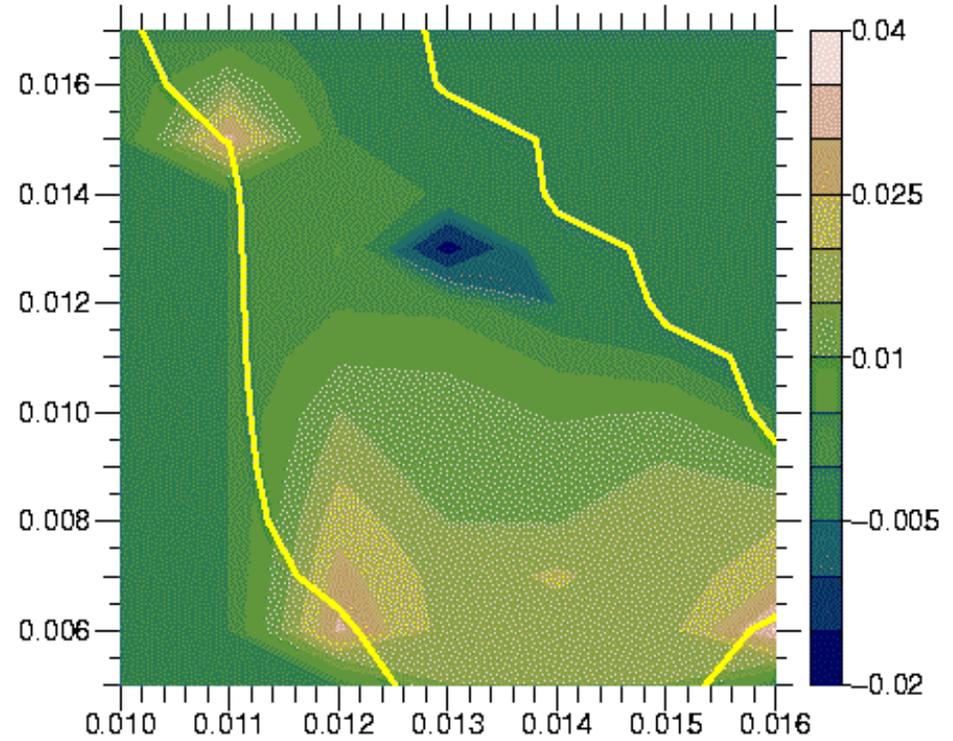
We can do the same optimum finding on the new space



Shift causing minimum fmha



Shift furthest from yield cliffs



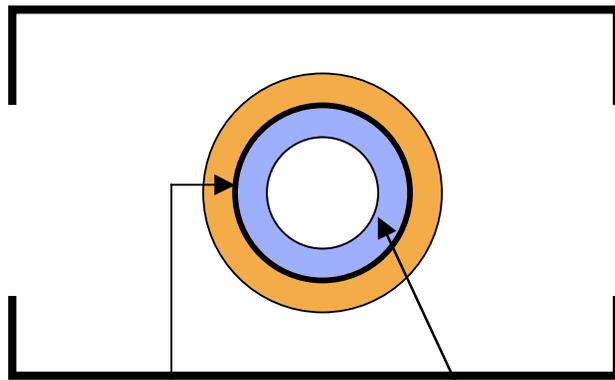
Ablator thickness (cm)

Again we find a shift of 100-200 ps to be optimum in the most interesting region

While tuning in spherical geometry is nice, experimental tuning is done in planar geometry



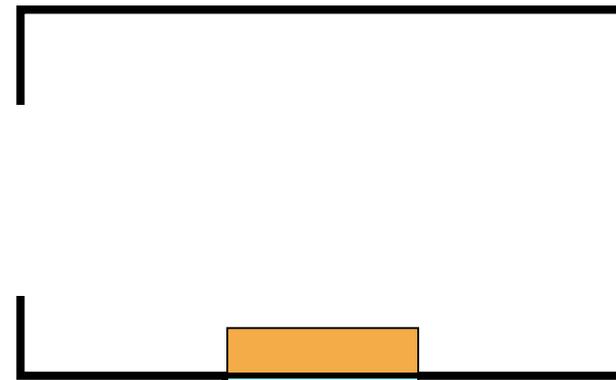
“Spherical Tuning”



Pressures measured on interface in spherical geometry

DT Ice

Experimental Tuning



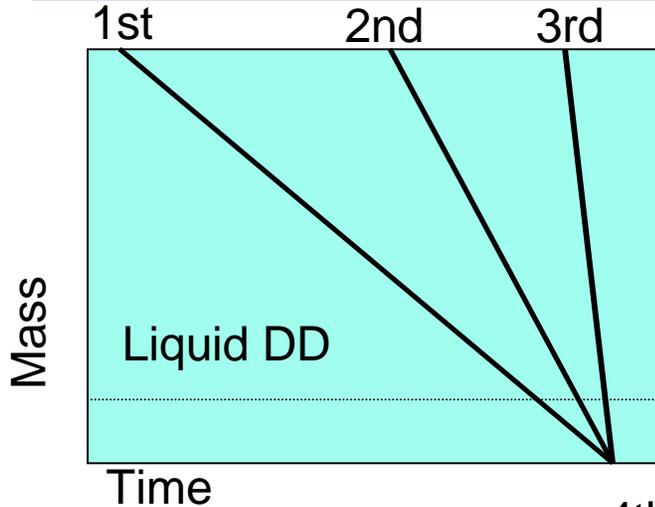
Pressures measured on interface in planar geometry

Liquid DD

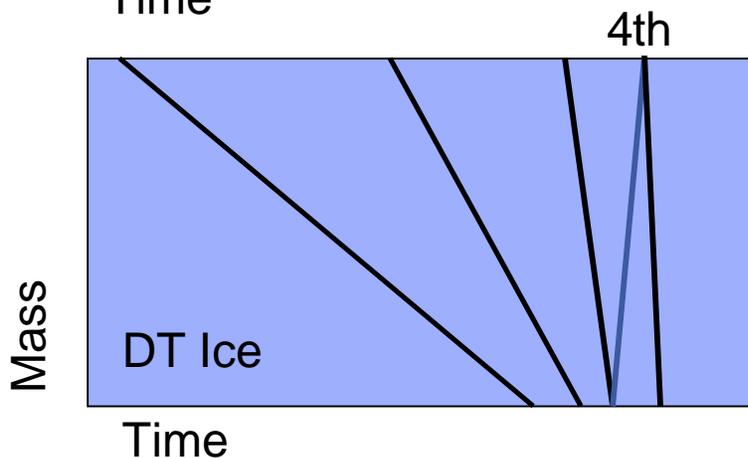
Dave Munro has shown that the same drive that leads to the proper coalescence depth in liquid DD in a planar geometry provides a good pulse shape for spherical capsules using DT ice.

How can we implement this pulse shaping technique?

Pulse shaping using a planar geometry is a two stage process



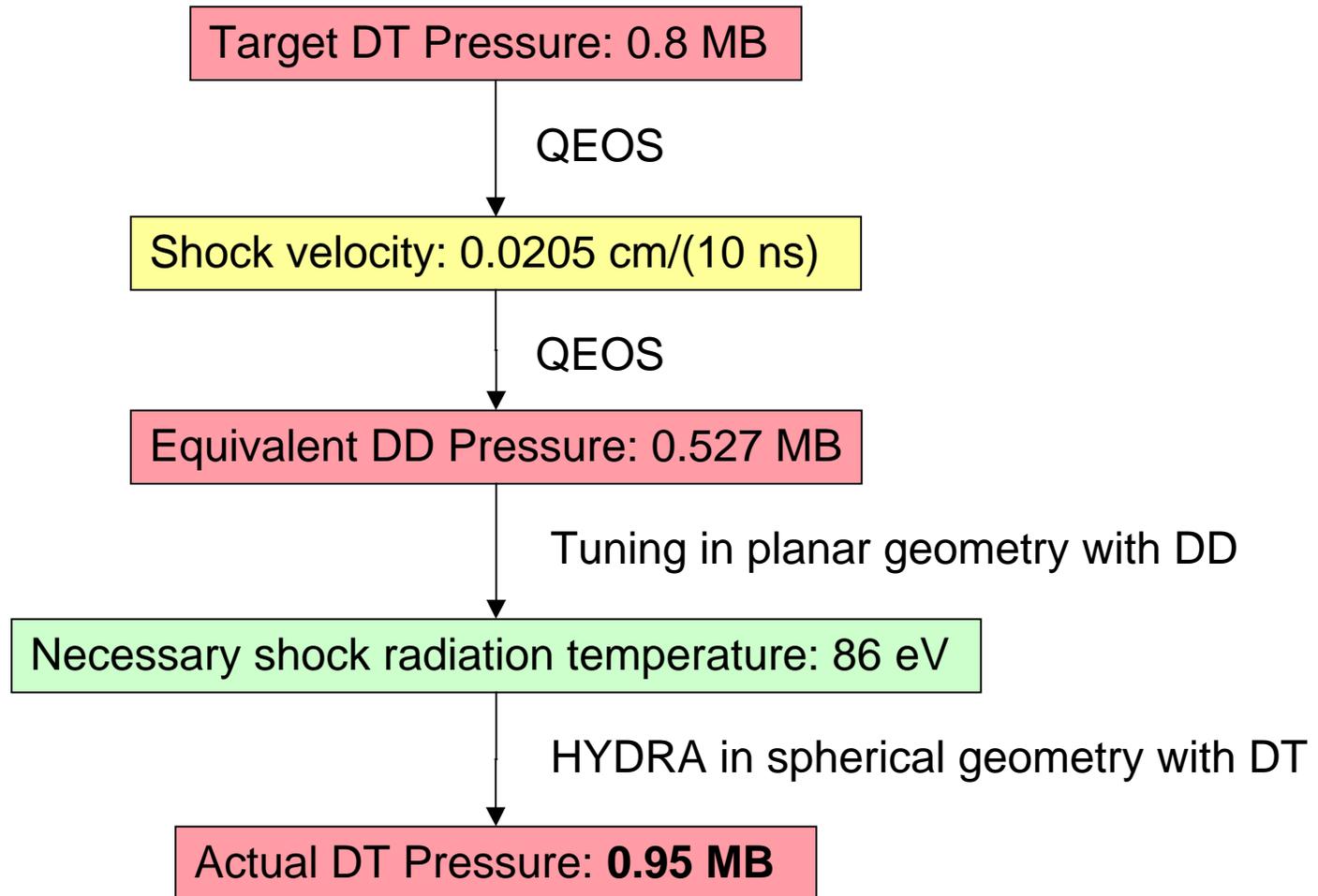
First, the first three shocks are tuned to get the proper convergence depth in the planar geometry with liquid DD



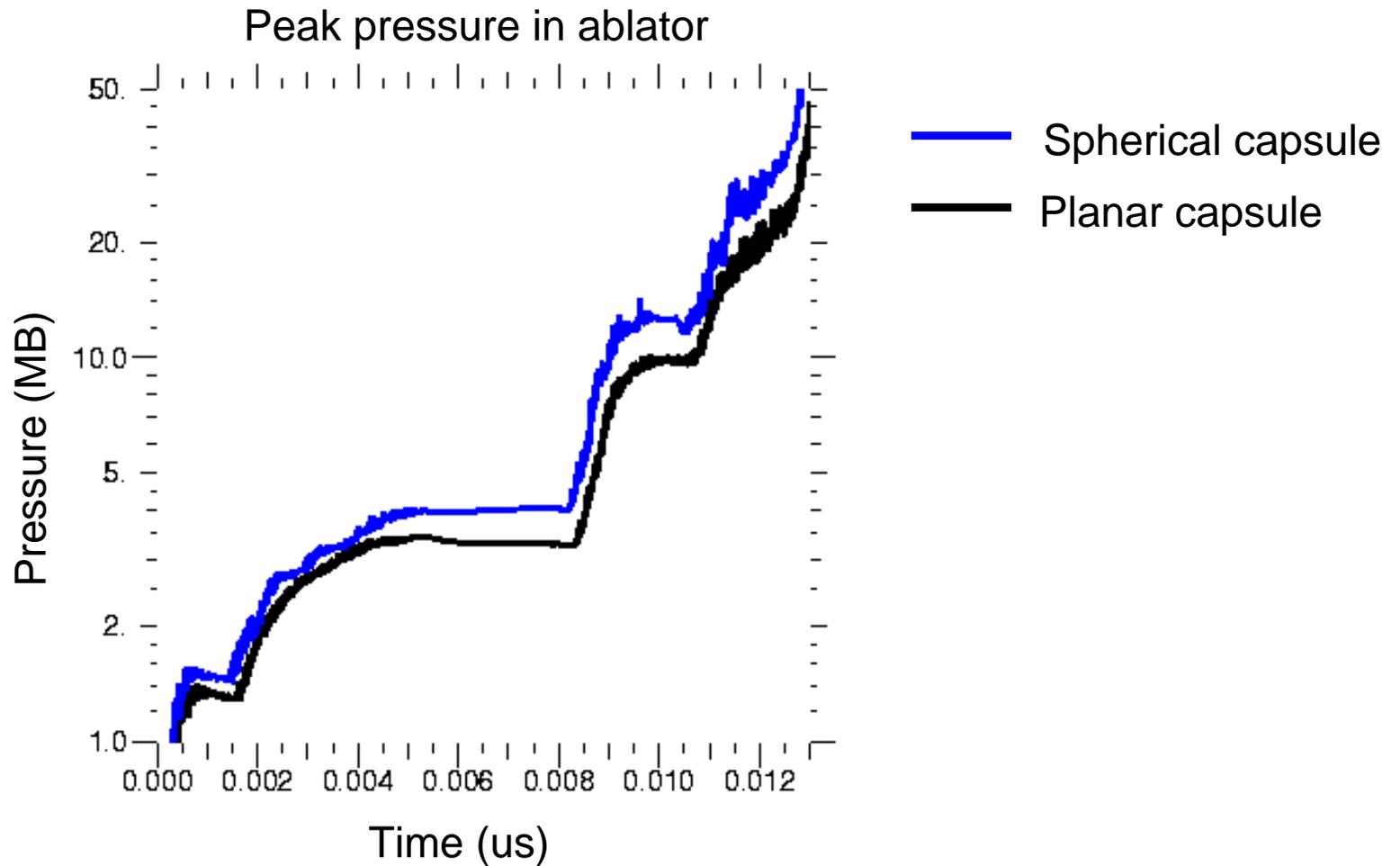
Then, the first three shocks are fixed and the fourth shock is timed in a spherical geometry with DT Ice.

While the shock velocities should be the same in both cases, the pressure at the interface is different due to the different material properties

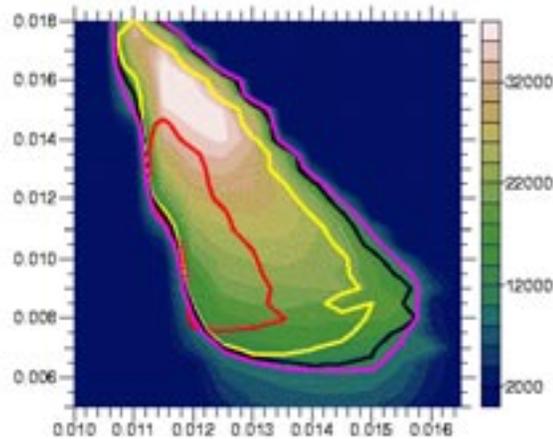
This process, however, does not yield desired results.



The culprit is the difference between planar and spherical ablation



Summary



- We have implemented an automated pulse shaper for NIF capsules in HYDRA.
- We have developed the infrastructure to do scans using the automatic pulse shaper across any n-dimensions of capsule parameter space.
- Using this infrastructure, we have performed several scans examining parameters for uniformly doped Beryllium capsules.
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Further Work



- **Work on solutions to pulse-shaping in planar geometry**
- **Using this pulse shaper, complete scans over other pulse-shaping parameters.**
- **Implement the ability to use 2D decks to determine robustness of capsules.**
- **Make scanning apparatus and pulse shaper completely compatible with graded dopants.**