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## Parallel auxiliary space AMG for definite Maxwell problems

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(joint work with Panayot S. Vassilevski)

Motivated by the needs of large multi-physics simulation codes, we are interested in algebraic solvers for the linear systems arising in time-domain electromagnetic simulations. Our focus is on finite element discretizations, and we are developing scalable parallel preconditioners which employ only fine-grid information, similar to algebraic multigrid (AMG) for diffusion problems. In the last few years, the search for efficient algebraic preconditioners for  $\mathbf{H}(\text{curl})$  bilinear forms has intensified. The attempts to directly construct AMG methods had some success, see [12, 1, 7]. Exploiting available multilevel methods on auxiliary mesh for the same bilinear form led to efficient auxiliary mesh preconditioners to unstructured problems as shown in [4, 8]. A computationally more attractive approach was recently proposed by Hiptmair and Xu [5]. In contrast to the auxiliary mesh idea, the method in [5] uses a nodal  $\mathbf{H}^1$ -conforming auxiliary space on the same mesh. This significantly simplifies the computation of the corresponding interpolation operator.

In the present talk, we consider several options for constructing unstructured mesh AMG preconditioners for  $\mathbf{H}(\text{curl})$  problems and report a summary of computational results from [10, 9]. Our approach is slightly different than the one from [5], since we apply AMG directly to variationally constructed coarse-grid operators, and therefore no additional Poisson matrices are needed on input. We also consider variable coefficient problems, including some that lead to a singular matrix. Both type of problems are of great practical importance and are not covered by the theory of [5].

We are interested in solving the following variational problem stemming from the definite Maxwell equations:

$$(1) \quad \text{Find } \mathbf{u} \in \mathbf{V}_h : \quad (\alpha \text{curl } \mathbf{u}, \text{curl } \mathbf{v}) + (\beta \mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}), \quad \text{for all } \mathbf{v} \in \mathbf{V}_h .$$

Here we consider  $\alpha > 0$  and  $\beta \geq 0$  which are scalar coefficients, but extensions to (semi)definite tensors are possible. We allow  $\beta$  to be zero in part or the whole domain (in which case the resulting matrix is only semidefinite and for solvability, the right-hand side should be chosen to satisfy compatibility conditions). Let  $\mathbf{A}_h$  be the stiffness matrix corresponding to (1), where  $\mathbf{V}_h$  is the (lowest order) Nédélec space associated with a triangulation  $\mathcal{T}_h$ .

Let  $S_h$  be the space of continuous piecewise linear finite elements associated with the same mesh  $\mathcal{T}_h$  as  $\mathbf{V}_h$ , and  $\mathbf{S}_h$  be its vector counterpart. Let  $G_h$  and  $\mathbf{I}_h$  be the matrix representations of the mapping  $\varphi \in S_h \mapsto \nabla \varphi \in \mathbf{V}_h$  and the nodal interpolation from  $\mathbf{S}_h$  to  $\mathbf{V}_h$ , respectively. Note that  $G_h$  has as many rows as the number of edges in the mesh, with each row having two nonzero entries:  $+1$  and  $-1$  in the columns corresponding to the edge vertices. The sign depends on the orientation of the edge. Furthermore,  $\mathbf{I}_h$  can be computed based only on  $G_h$  and on the coordinates of the vertices of the mesh.

The auxiliary space AMG preconditioner for  $\mathbf{A}_h$  is a subspace correction method utilizing the subspaces  $\mathbf{V}_h$ ,  $G_h\mathbf{S}_h$ , and  $\mathbf{I}_h\mathbf{S}_h$ . Its additive form reads (cf. [13])

$$(2) \quad A_h^{-1} + G_h B_h^{-1} G_h^T + \mathbf{I}_h \mathbf{B}_h^{-1} \mathbf{I}_h^T,$$

where  $A_h$  is a smoother for  $\mathbf{A}_h$ , while  $B_h$  and  $\mathbf{B}_h$  are efficient preconditioners for  $G_h^T \mathbf{A}_h G_h$  and  $\mathbf{I}_h^T \mathbf{A}_h \mathbf{I}_h$  respectively. Since these matrices come from elliptic forms, the preconditioner of choice is AMG (especially for unstructured meshes).

If  $\beta$  is identically zero, one can skip the subspace correction associated with  $G_h$ , in which case we get a two-level method.

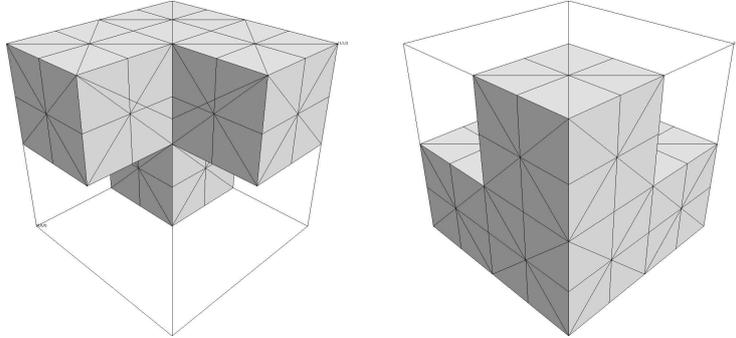
The motivation for (2) is that any finite element function  $\mathbf{u}_h \in \mathbf{V}_h$  allows for decomposition of the form (cf., [5])  $\mathbf{u}_h = \mathbf{v}_h + \mathbf{I}_h \mathbf{z}_h + \nabla \varphi_h$  with  $\mathbf{v}_h \in \mathbf{V}_h$ ,  $\mathbf{z}_h \in \mathbf{S}_h$  and  $\varphi_h \in S_h$  such that the following stability estimates hold,

$$(3) \quad h^{-1} \|\mathbf{v}_h\|_0 + \|\mathbf{z}_h\|_1 \leq C \|\operatorname{curl} \mathbf{u}_h\|_0 \quad \text{and} \quad \|\nabla \varphi_h\|_0 \leq C \|\mathbf{u}_h\|_0.$$

A parallel solution algorithm based on (2) was implemented in the *hypr* library [6], under the name AMS (Auxiliary space Maxwell Solver). The internal AMG V-cycles employ *hypr*'s algebraic multigrid solver BoomerAMG [2]. Our results on unstructured meshes in two and three dimensions, including problems with variable coefficients and zero conductivity, clearly demonstrate the scalability of this preconditioner on hundreds of processors. A sample of the numerical experiments to be presented is shown in Table 1, using the following notation:  $np$  is the number of processors in the run,  $N$  is the global size of the problem, and  $t$  denotes the average time to solution (in seconds) on a machine with 2.4GHz Xeon processors.

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$np$	$N$	$p$									$t$
		-8	-4	-2	-1	0	1	2	4	8	
1	83,278	9	9	9	9	9	9	10	11	11	5s
2	161,056	10	10	10	10	10	10	10	11	11	9s
4	296,032	11	12	12	12	11	11	12	13	13	9s
8	622,030	13	13	13	12	12	12	13	15	14	12s
16	1,249,272	13	13	13	13	13	13	13	15	14	13s
32	2,330,816	15	15	15	15	15	15	15	16	15	14s
64	4,810,140	16	16	16	16	16	15	16	18	17	17s
128	9,710,856	16	16	16	16	16	16	16	17	17	23s
256	18,497,920	19	19	19	19	19	19	19	21	20	27s
512	37,864,880	21	20	20	20	20	20	20	23	22	32s
1024	76,343,920	20	20	20	20	20	20	20	21	21	56s

TABLE 1. Number of AMS-PCG iterations with tolerance  $10^{-6}$ , for the problem (1) on the unit cube with  $\alpha = 1$ , and  $\beta \in \{1, 10^p\}$  having different values in the shown regions (cf. [3]).

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